

MATERIALS FOR ENERGY HARVESTING

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ENERGY HARVESTING

System at
non –equilibrium



System at
equilibrium

Water mill



Automatic watch



(Electric) Energy

SOURCES

► Thermal gradient

- Thermoelectric materials (semiconductors, clathrates, skutterudite)



► Motion/mechanical vibrations

- Piezoelectrics, magnets, electrets, magnetostrictive



► Photovoltaics/EM background “noise”

- Si PV cell, rectenna, chem. harvester



SUMMARY

▶ Introduction

- ▶ Materials and devices

▶ Thermoelectrics

- ▶ Physical properties
- ▶ Phononic crystals

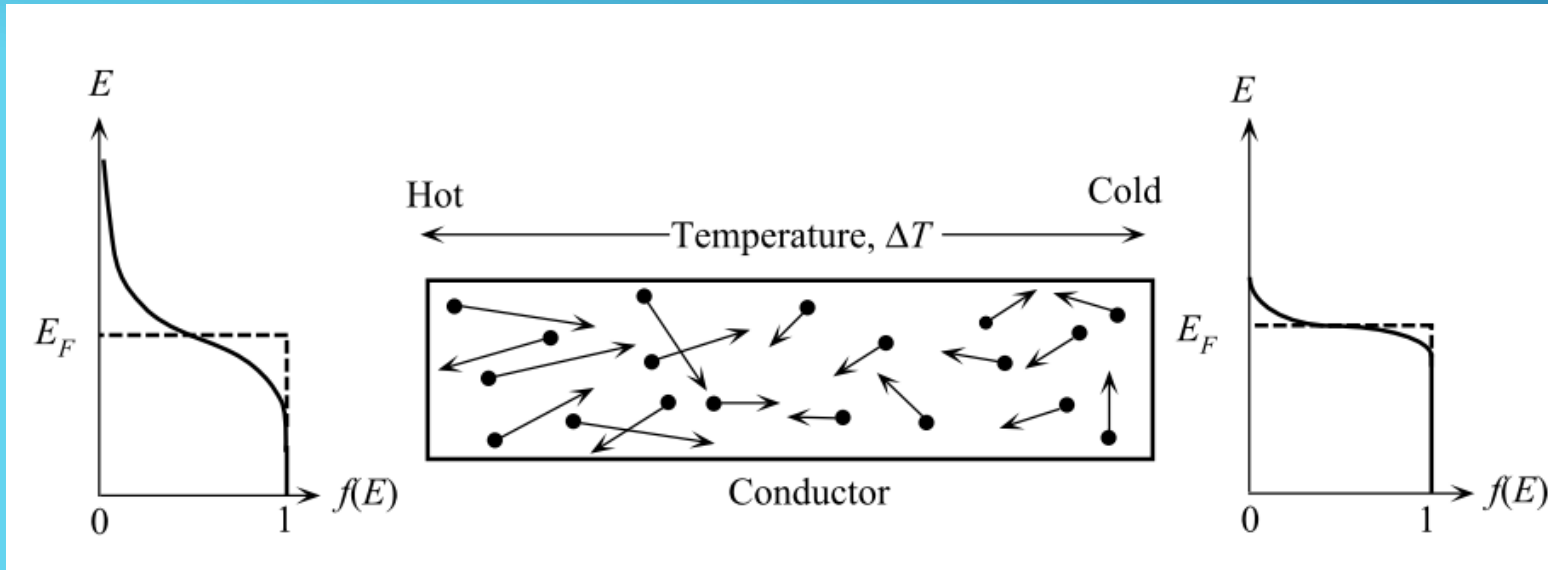
▶ Piezoelectrics

- ▶ Physical properties/fabrication
- ▶ ZnO microrods

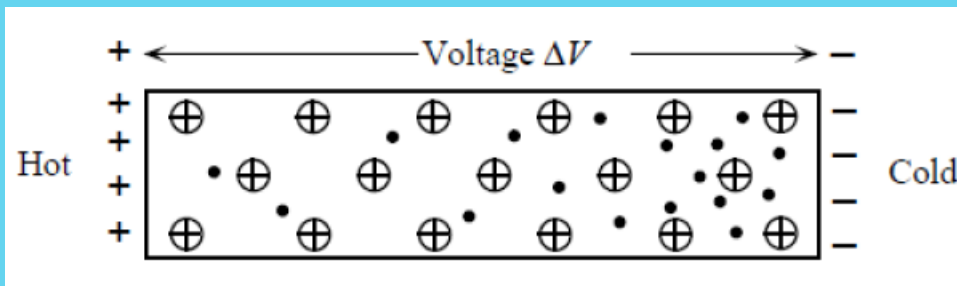
▶ Electrets

- ▶ Physical properties/fabrication
- ▶ SiO₂ micro particles as electrets

THERMOELECTRIC EFFECT



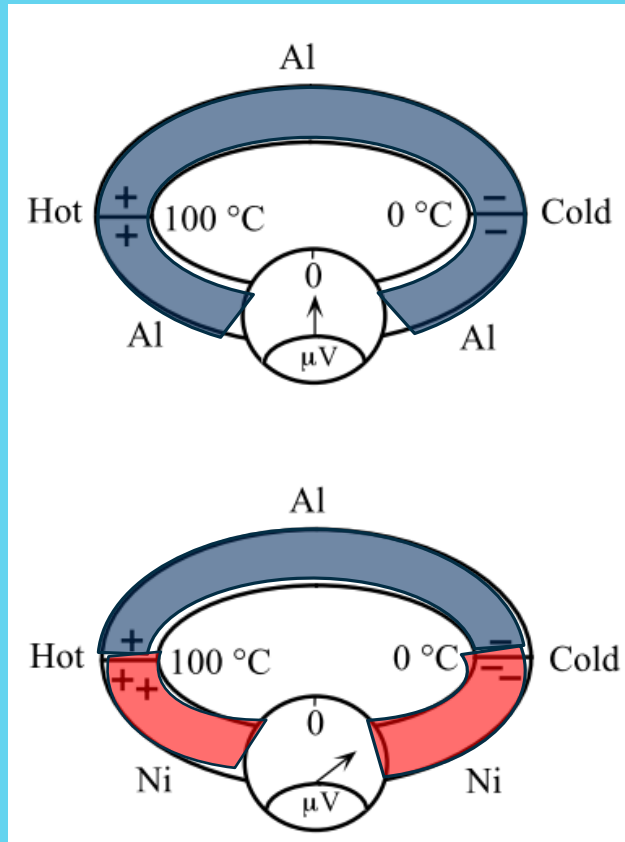
A temperature difference between two points in a conductor gives rise to a voltage difference



Seebeck coefficient

$$S = \frac{\Delta V}{\Delta T}$$

THERMOCOUPLES



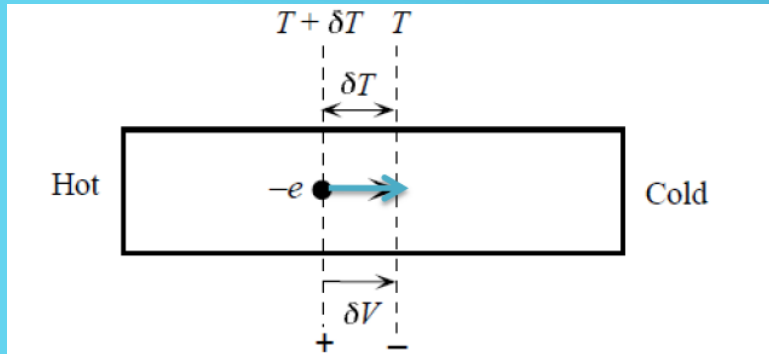
This effect is exploited in thermocouples.

In a circuit made by a single wire with cold and hot regions no voltage difference is observed.

However if we have the junction between different metals, as the Seebeck effect is different in different a net voltage can be observed.

As for a fixed couple of metal, the value depends only on the temperature difference, this device can be used as a sensor

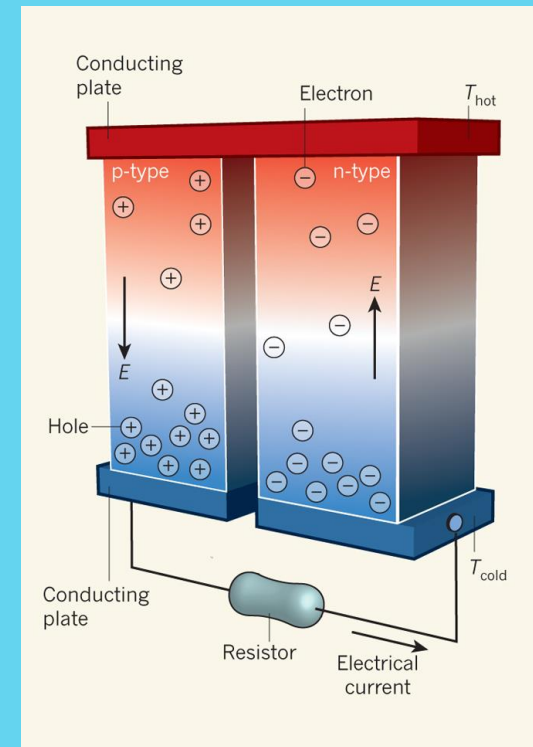
THERMOELECTRIC GENERATOR



The «hot» electrons, faster, have higher possibility to move toward the cold region than the «cold» electrons in the opposite direction.

This leads to an increase of the electric potential energy which can be exploited

Just like for the thermocouple, we need to close the circuit maintaining the asymmetry in the charge carrier behavior

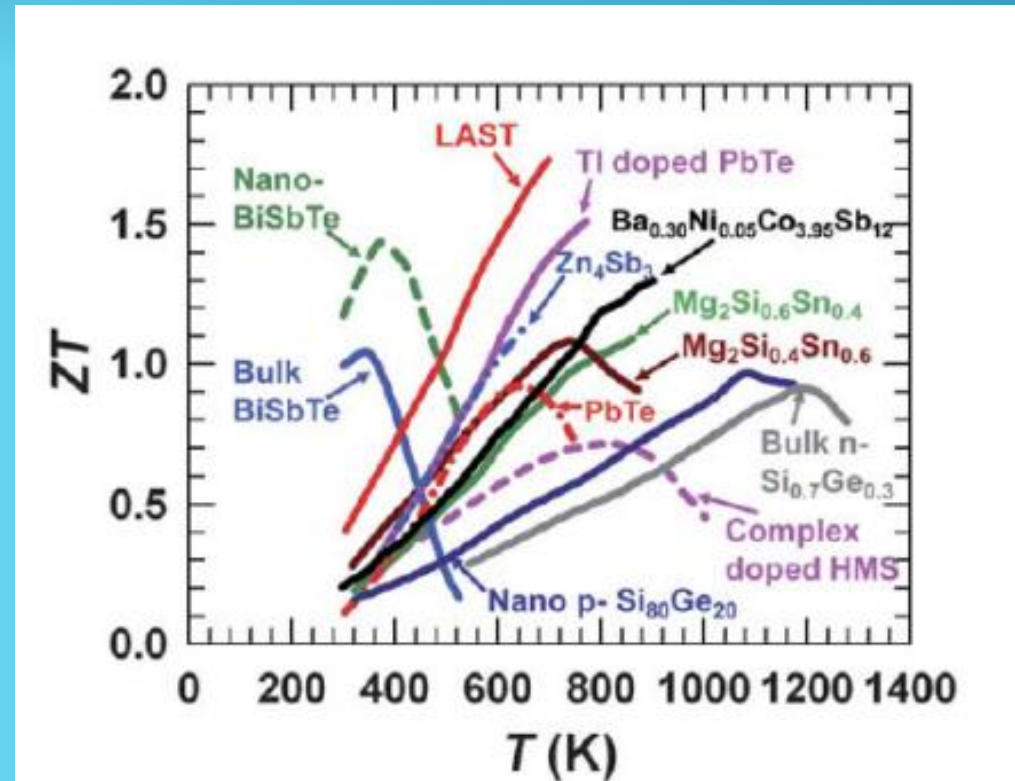


COMPARISON OF THERMOELECTRIC MATERIALS

Electric conductivity

$$ZT = \frac{\sigma_e S^2}{\kappa_L} T$$

Thermal conductivity

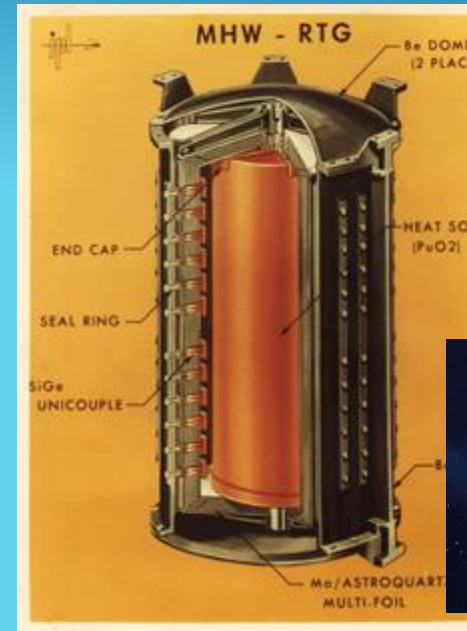
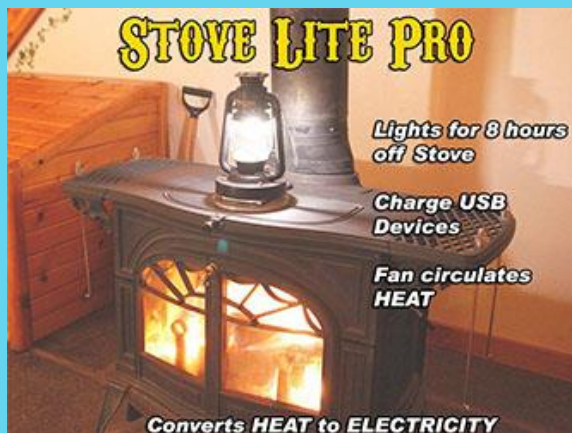


Phonon glass and electron crystal

SOME DEVICES



GENTHERM
GLOBAL POWER TECHNOLOGIES

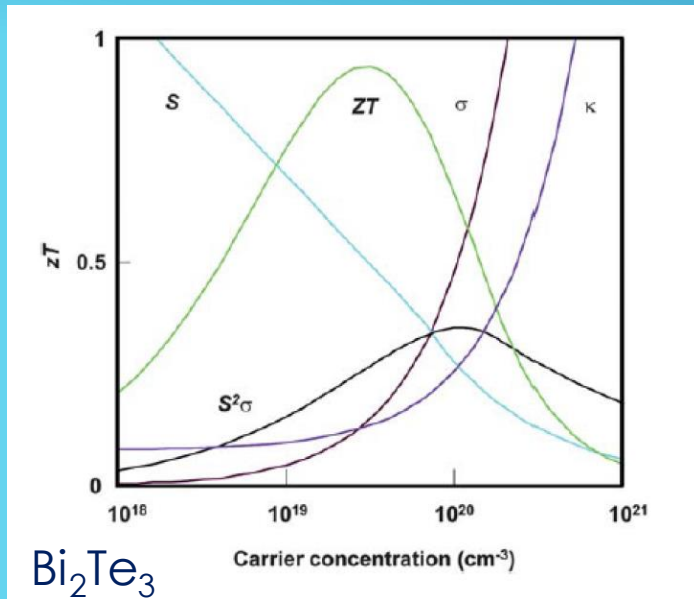


Radioisotope
thermoelectric generator

Low efficiency, (3-4%, for $ZT = 1-2$)

but highly reliable (No moving parts)

MATERIAL OPTIMIZATION



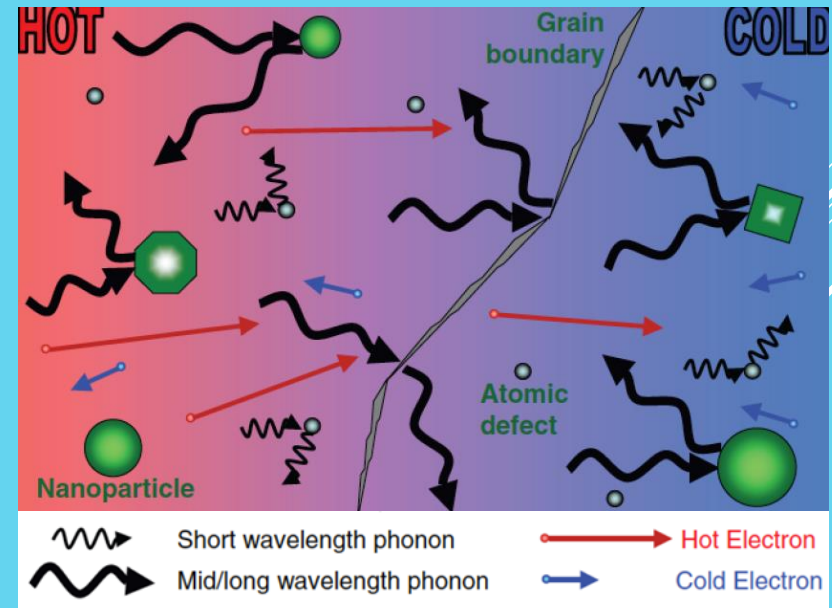
$$ZT = \frac{\sigma_e S^2}{\kappa_L} T$$

$$\kappa_L = \kappa_e + \kappa_p$$

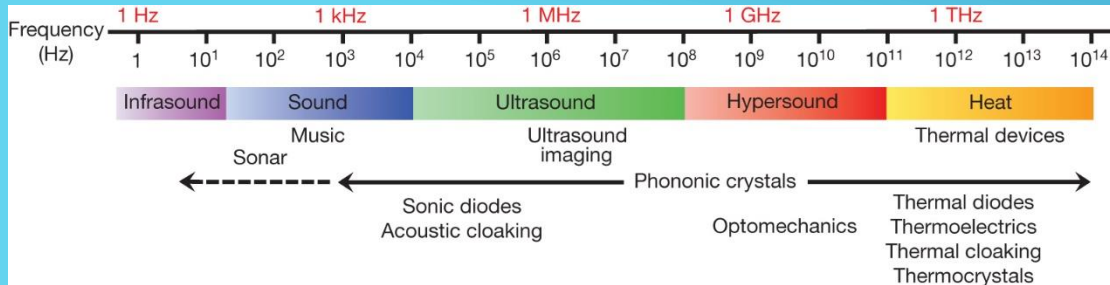
Electric conductivity and electronic part of the thermal conductivity are strictly related.

Phonons are more sensitive than electrons to discontinuities in the materials.

Defects and interfaces can decrease the phononic thermal conductivity



SOUND AND HEAT REVOLUTIONS IN PHONONICS

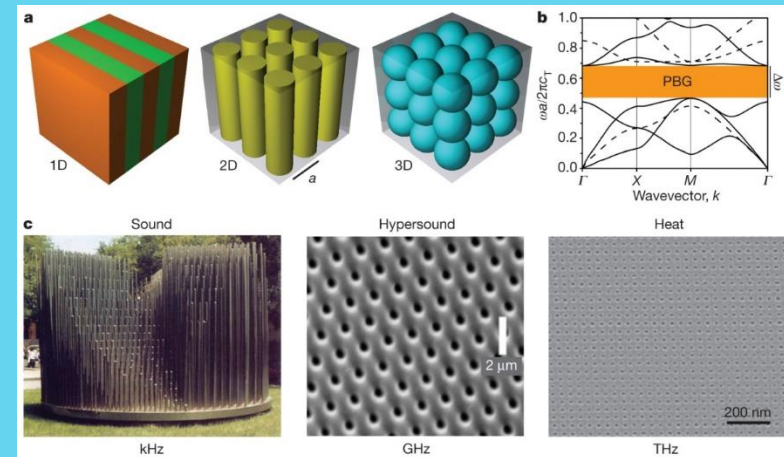


The phononic spectrum

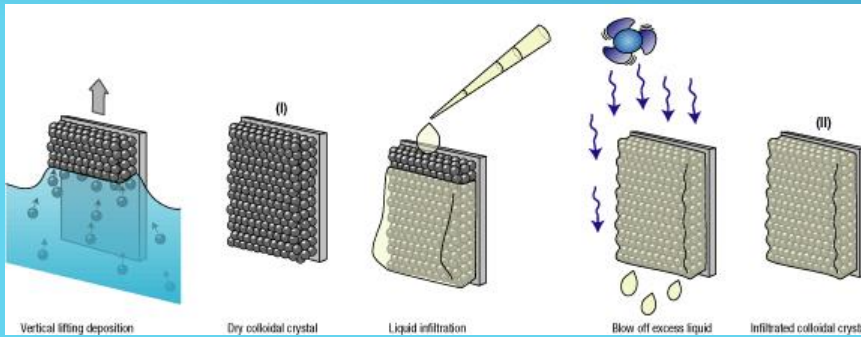
The development of sonic and thermal devices requires the design, fabrication and characterization of composite materials ranging from the centimetre scale to the nanometre scale

In phononic crystals, mechanical waves with frequencies within a specific range are not allowed to propagate within the periodic structure. This '**phononic bandgap**' owing to wave interference effects, occurs for phonon wavelengths that are comparable to the structure periodicity.

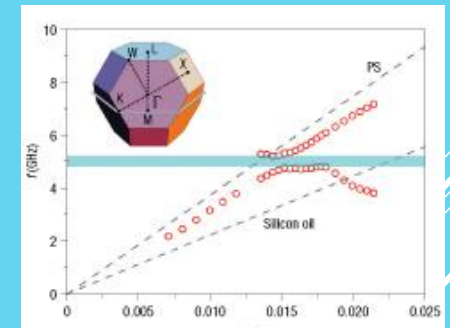
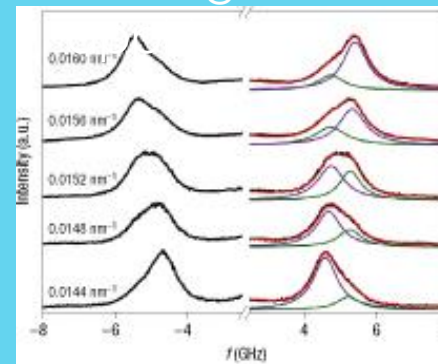
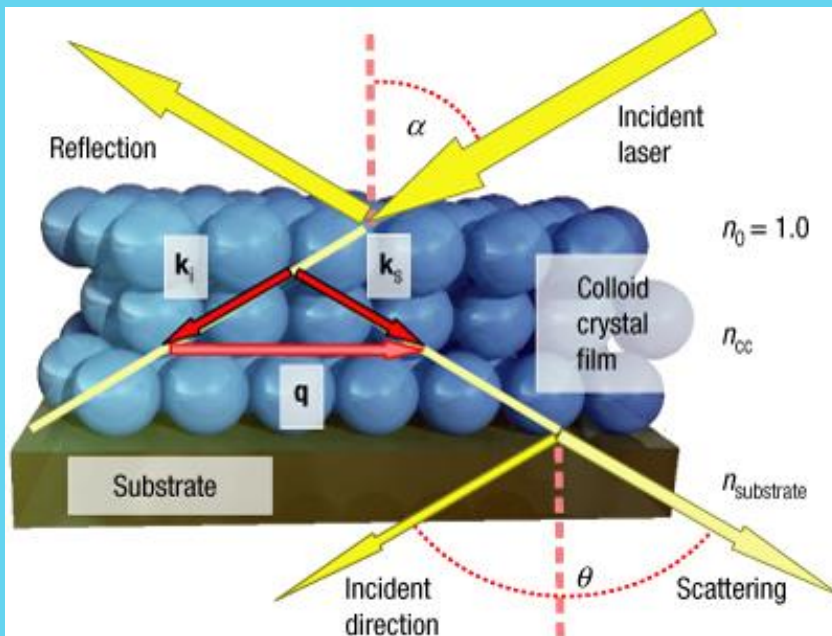
Phononic crystals



COLLOIDAL CRYSTALS

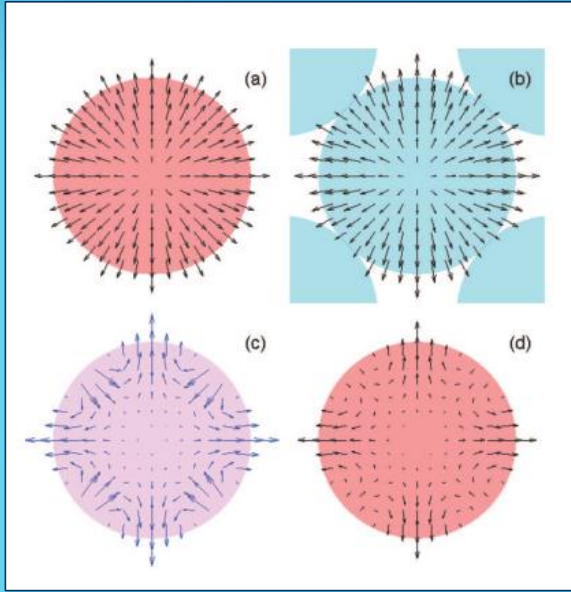


An ordered array of colloid particles (PS, 100-1000 nm) fabricated by self-assembly in a fcc structure. It presents phononic band gap, making it valuable for devices exploiting phonon control (acoustic isolator, heat management, acousto-optical



A matching index liquid infiltration is needed to overcome multiple scattering (and loss of q determination). But there is also an acoustic impedance matching.

CLUSTER AND CRYSTAL MODES



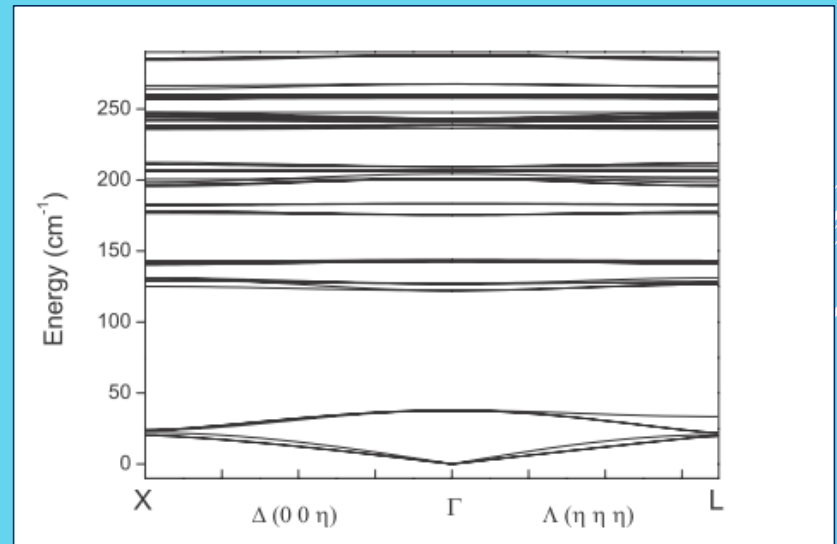
Band arise from the interaction between near particles

Coupling between different cluster modes, but limited lifting of the degeneracy.

Band Gaps in Phononic Crystals:

1) A mesoscopic periodicity breaks continuous acoustic bands

2) A tight binding approach: the weak interaction among the particles will couple the vibrations with the result that the discrete vibrational spectrum ω_α ($\alpha = (\text{pnlm})$) will be transformed into vibrational bands.



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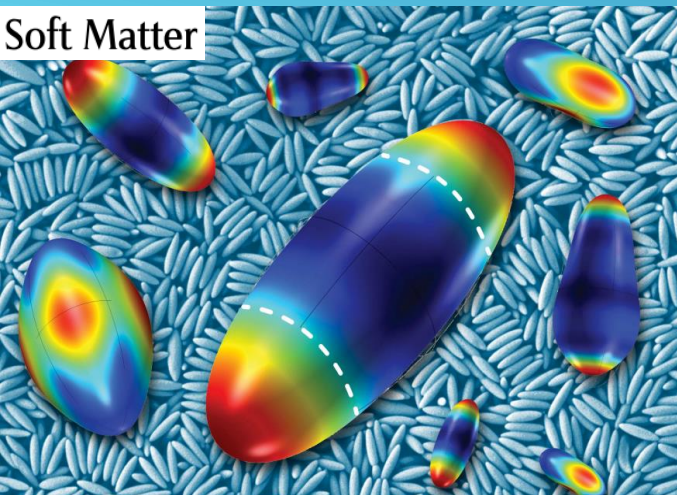
Phononic crystals of spherical particles: A tight binding approach

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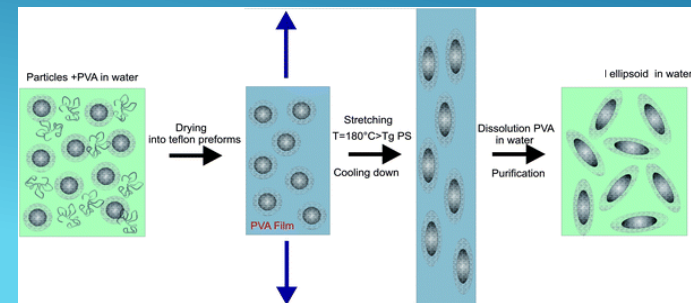


Soft Matter

Elongated polystyrene spheres as resonant building blocks in anisotropic colloidal crystals†

Dirk Schneider,^a Peter J. Beltramo,^b Maurizio Mattarelli,^c Patrick Pfeleiderer,^d Jan Vermant,^d Daniel Crespy,^a Maurizio Montagna,^c Eric M. Furst^b and George Fytas^{*ae}

Polystyrene particles were stretched by a matrix assisted elongation process



➤ Elastic Properties

➤ Normal modes

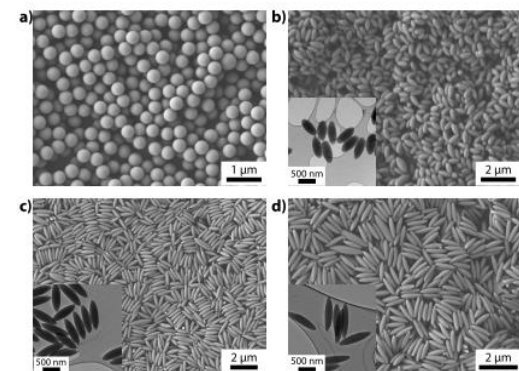


Fig. 1 SEM micrographs of seed spheres S0 (a) and spheroids S2–S4 (b–d) of various aspect ratio. Insets denote TEM images of the respective particles.

Table 1 Size characterization of particles studied and material parameters used in calculations. a and b being the long and short axes of the spheroids, respectively. ρ is the mass density and n_{ref} the refractive index at 532 nm

ID	a/b	a/nm	b/nm	$\rho/\text{kg m}^{-3}$	n_{ref}^b
S0	1	400 ± 14	—	1050	1.599
S1 (ref. 30)	1.28 ± 0.07	459 ± 15	358 ± 16	965.1^a	1.540
S2	2.12 ± 0.16	690 ± 42	325 ± 15	921.9^a	1.511
S3	3.52 ± 0.20	986 ± 15	280 ± 15	869.4^a	1.477
S4	3.99 ± 0.22	1078 ± 93	270 ± 23	854.7^a	1.468

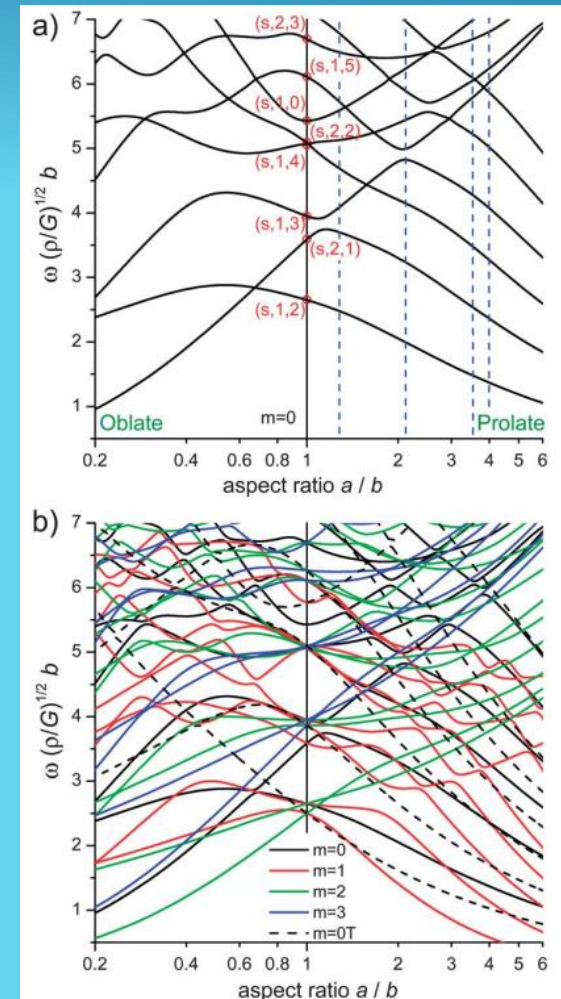
^a Calculated based on stretched particle volume. ^b Calculated from density *via* the Clausius–Mossotti relation.

SPHEROIDAL PARTICLES

- ▶ Rayleigh- Ritz method
axis ratio and sound velocities as
free parameters

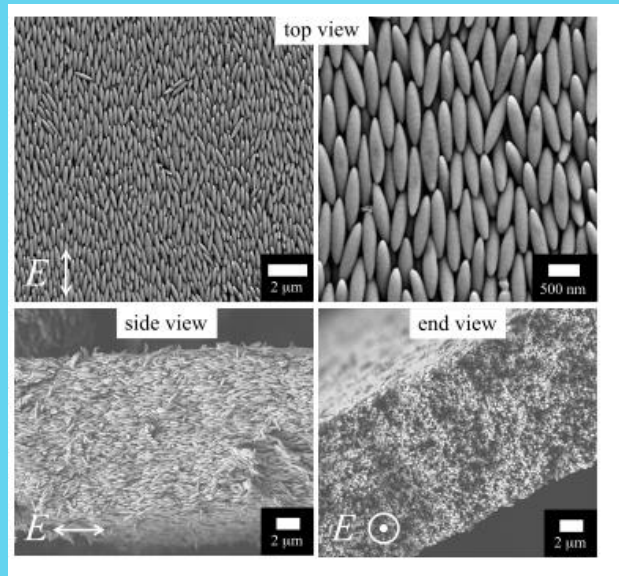
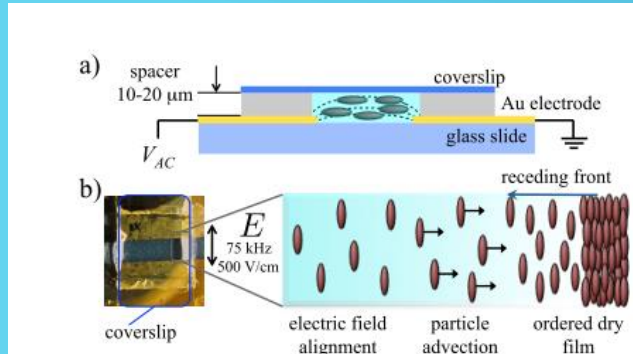


- ▶ Lifting of spherical degeneracy
(toward disk or beam modes)
- ▶ There is important mixing
between the various modes of
the sphere, but m and parity
remain good quantum numbers.

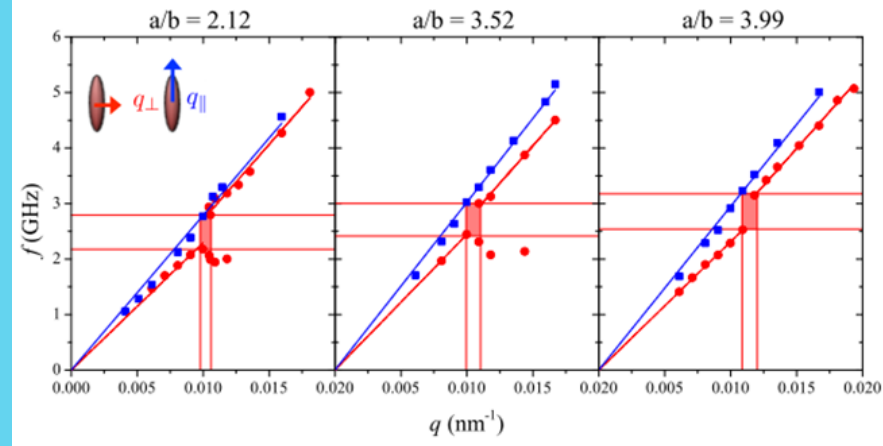
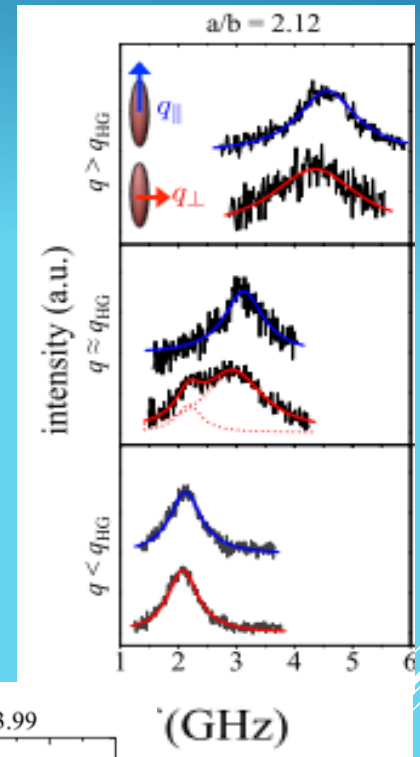


AXIS RATIO DISPERSION

ORIENTED ELONGATED PARTICLES

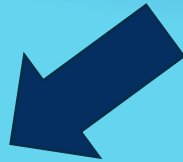


- The interaction of the acoustic phonon with other “optical” phonons depends on their symmetry
- Different gaps can occur for different propagation directions



TRANSDUCTION MECHANISMS AND MATERIALS

Mechanical action



Strain (cantilever, etc)



Motion («free» inertial mass,.)

Strain conversion

- ▶ Piezoelectrics
- ▶ Electroactive polymers
- ▶ Magnetostrictive mat.
- ▶ Magnetoelectrics

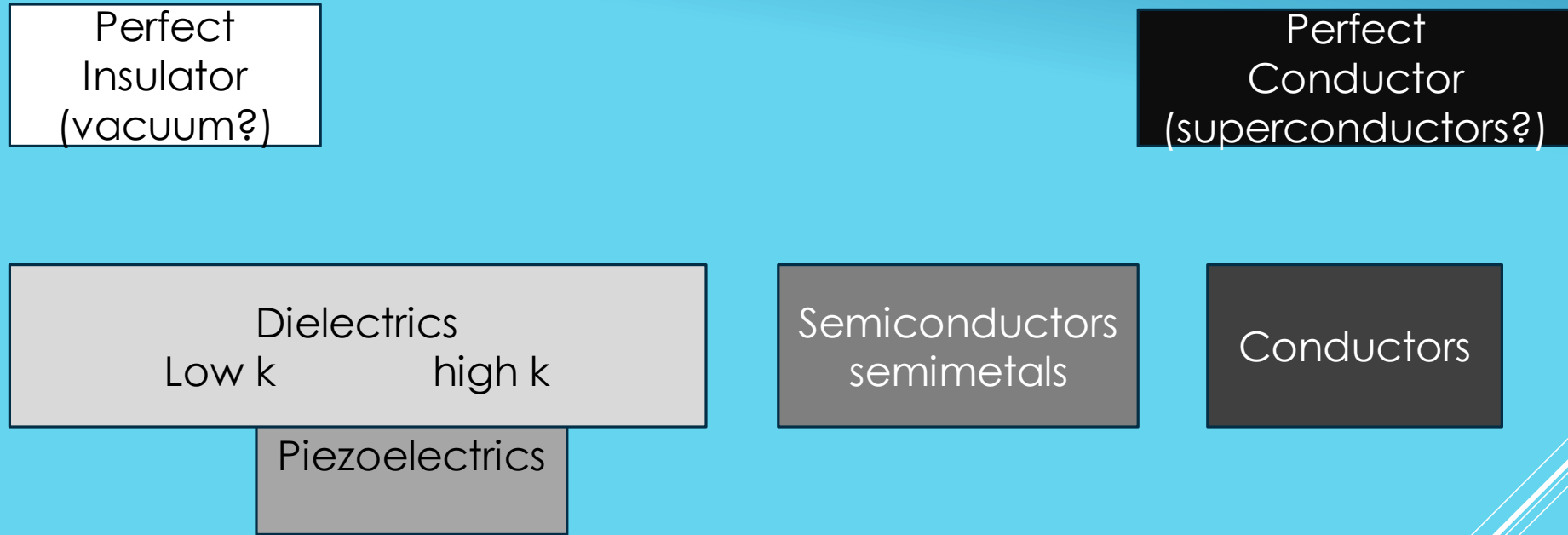
Electrostatic

- ▶ Electrets

Electromagnetic
Induction

- ▶ Magnets

ELECTRIC BEHAVIOR OF MATERIALS



Depending on the external conditions (electric field intensity or frequency, temperature, shape, strain) the behaviour of real materials can move between these extrema

ELECTRIC POLARIZATION

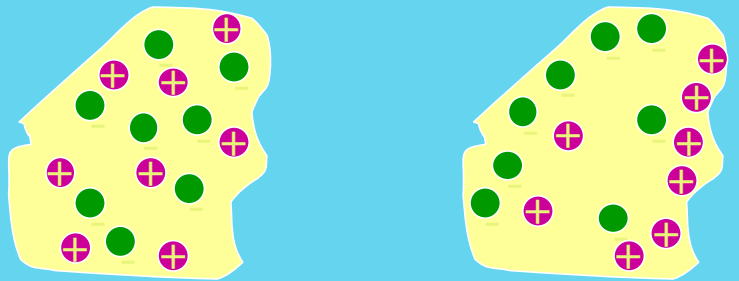
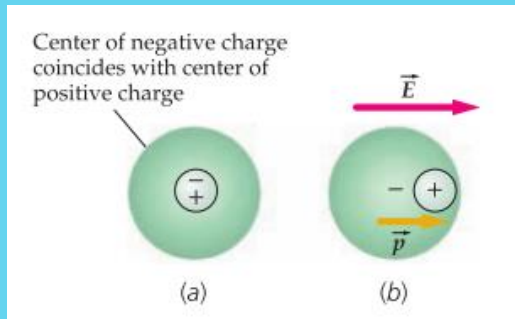
$$\mathbf{P} = \frac{\epsilon - 1}{4\pi} \mathbf{E}$$

In an external electric field, materials acquire an induced dipole moment or POLARIZATION

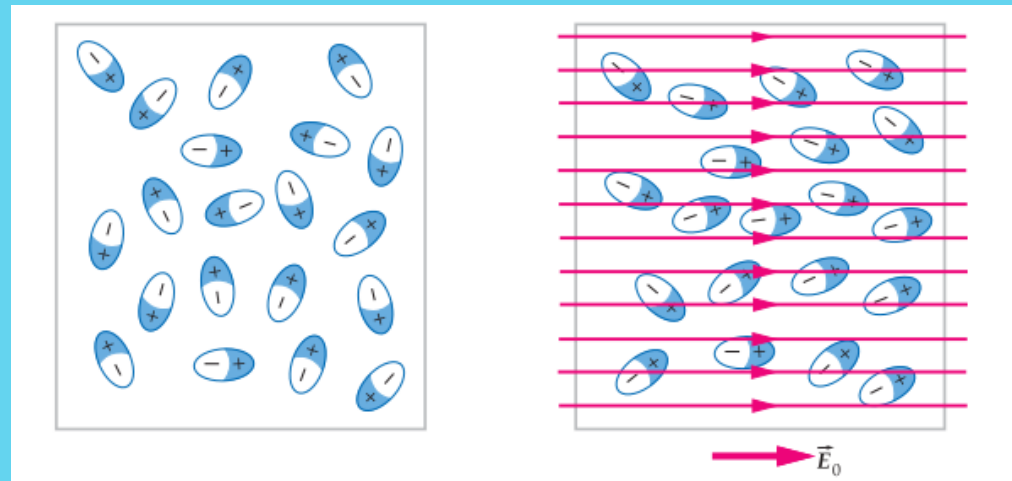
It can have different origins

Deformation

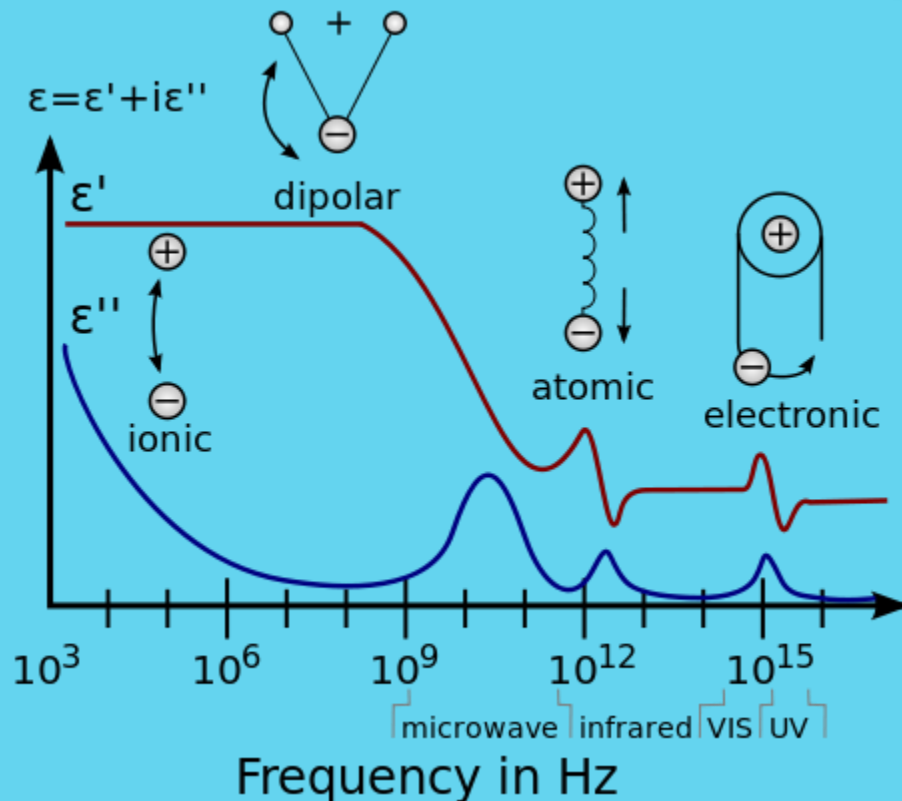
(of electron density or ions)



Orientation (of pre-existing dipoles)



FREQUENCY DEPENDENCE



- ▶ The response of the electroactive materials is strongly frequency dependent.
- ▶ They work best at resonance when the transferred power is maximum.
- ▶ They still react at lower frequency, while at higher frequency they cannot rearrange following the external field

MATERIALS IN MAGNETIC OR ELECTRIC FIELDS

Diamagnetic

Its atoms have no permanent magnetic moment. The induced moment opposes the external field.

Paramagnetic

Its atoms have permanent magnetic moment, which aligns with the external field

Ferromagnetic

Its atoms have permanent magnetic moment and strong exchange energy connects them into domains

Paraelectric

The induced dipole in dielectrics is always aligned with the external field

Ferroelectric

HYSTERESIS

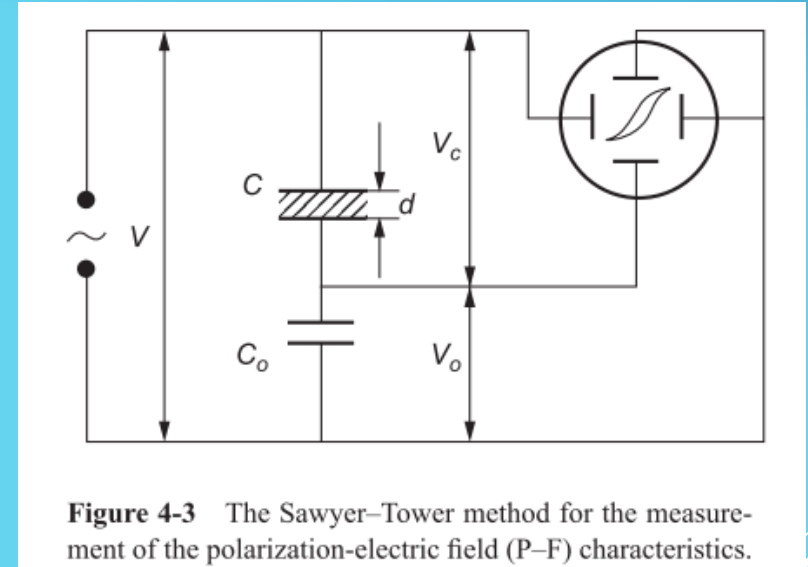
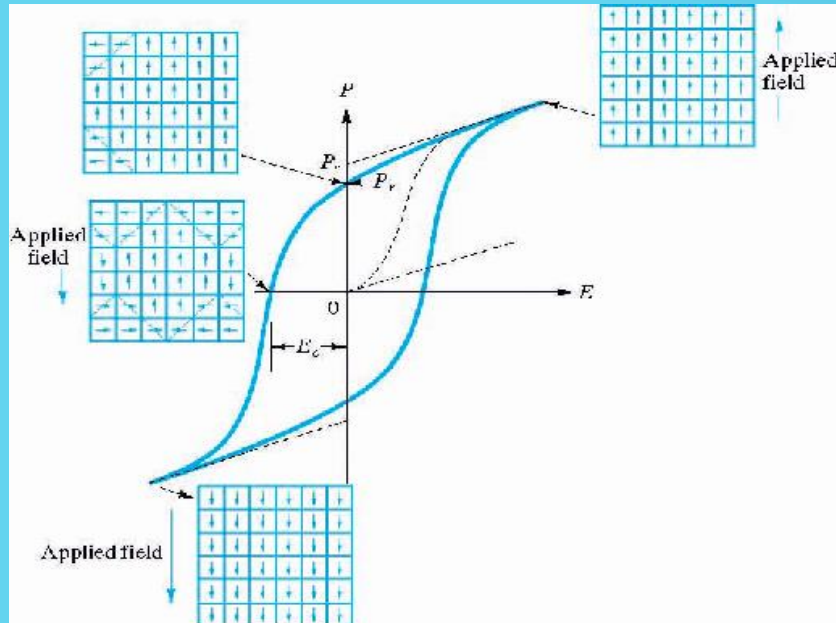


Figure 4-3 The Sawyer-Tower method for the measurement of the polarization-electric field (P - E) characteristics.

In Ferroelectrics (Ferromagnetics) materials, the thermal agitation cannot overcome the alignment of the domains. A coercive field is needed to reverse polarization (magnetization).

THE MICROSCOPIC ORIGIN

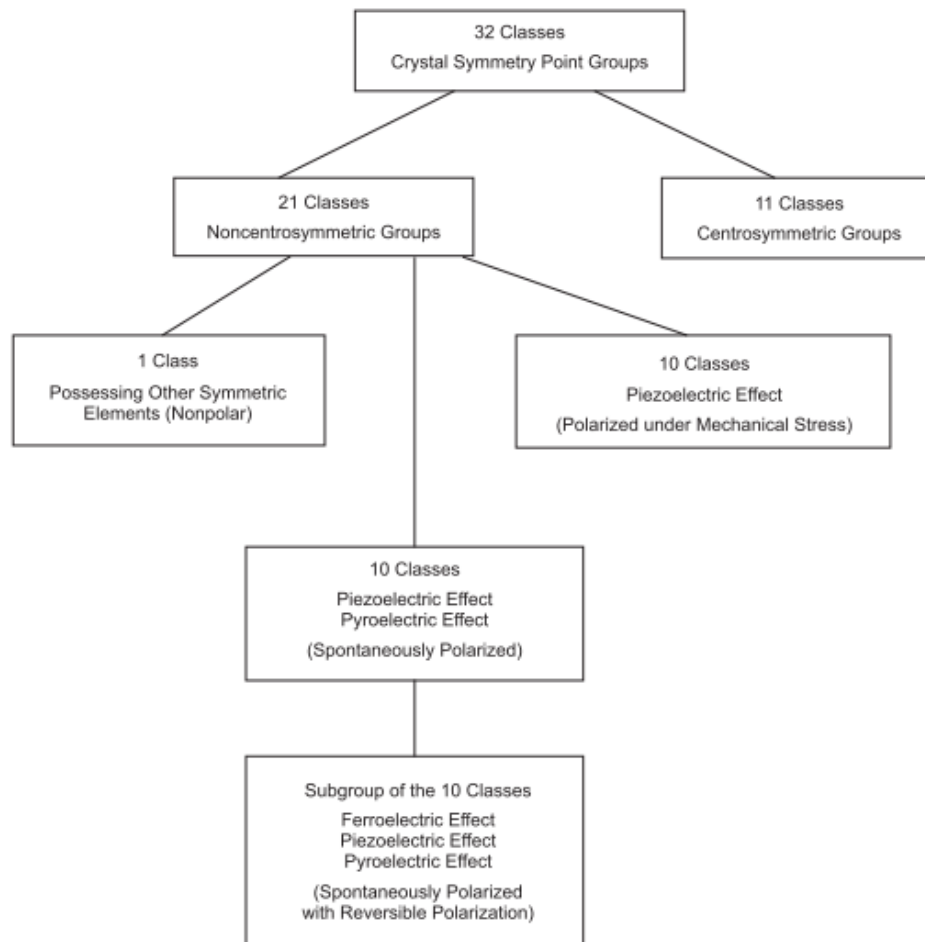
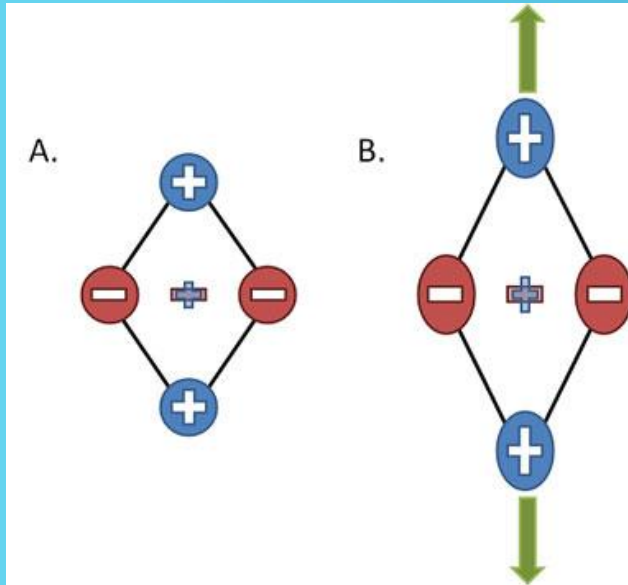


Figure 4-1 Classification of crystals showing the classes with piezoelectric, pyroelectric, and ferroelectric effects.

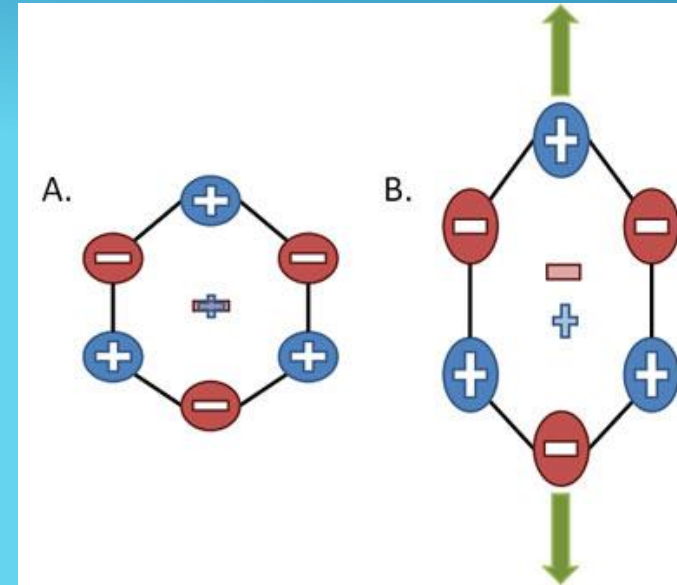
Symmetry of the unit cell is a necessary, but not sufficient condition.

Different charge distributions can cancel or strongly decrease the piezoelectric effect

PIEZOELECTRICITY



Centro-symmetric



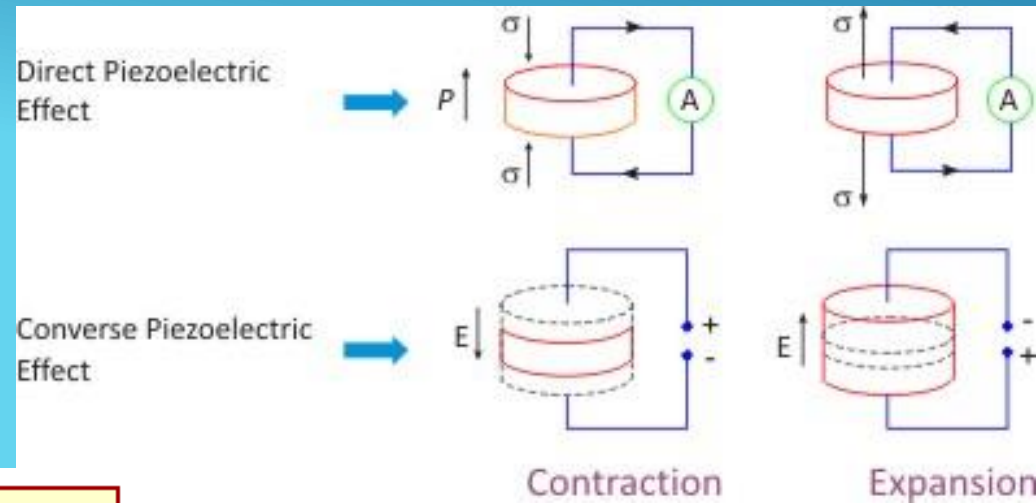
Non Centro-symmetric

In centrosymmetric crystals the strain does not move the center of charge of the positive or negative charges. On the other hand, in non centro-symmetric crystals, if the atoms have different charges because of the strain provokes the formation of an electric dipole.

PIEZOELECTRIC COEFFICIENTS

► $P = d \sigma$

► $\varepsilon = d^t E$



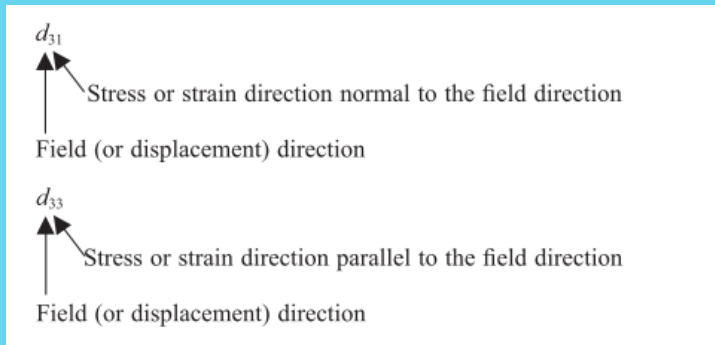
Material	Piezoelectric Constant, d (pm/V)
Quartz	2.3
Barium Titanate	100-149
Lead Niobate	80-85
Lead zirconate titanate	250-365

The direct piezoelectric effect is used as the basis for ENERGY HARVESTING (and force, pressure, vibration and acceleration sensors) while converse effect is used as a basis for actuator and displacement devices.

PIEZOELECTRIC COEFFICIENTS

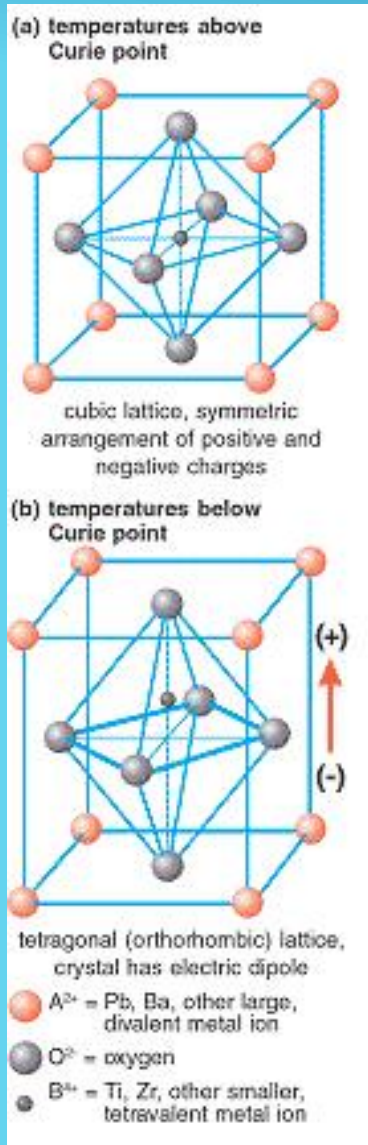
d is third-rank tensor (d_{ijk} 3x3x3), as it links the effect of strain/stress (second order tensor) to the induced Polarization/Electric field (vector).

However, it is often written in a contract matrix form (3x6), where 4,5,6 index are used to express shear stress/strain



Axis «3» is usually the anomalous axis of uniaxial piezoelectric crystals and the one where the effect is stronger

FERROELECTRICITY

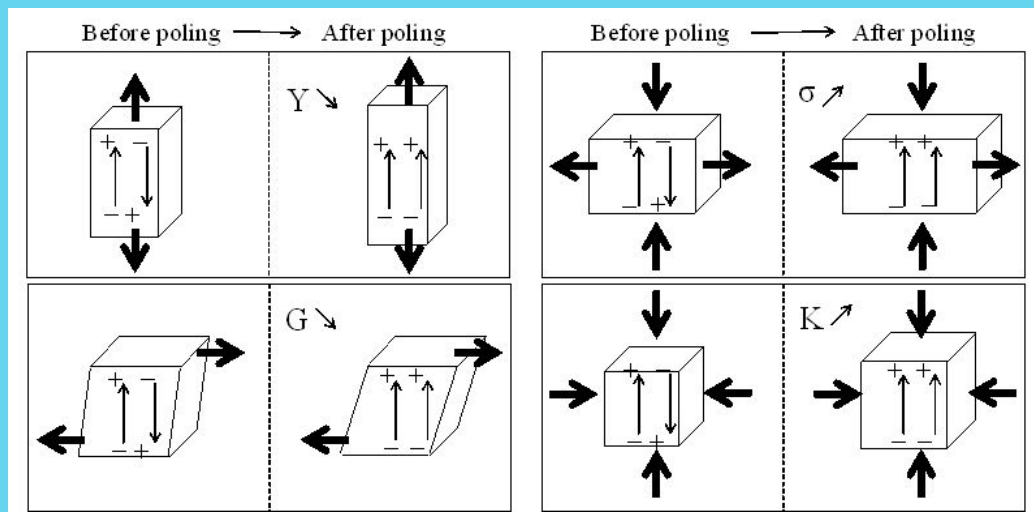
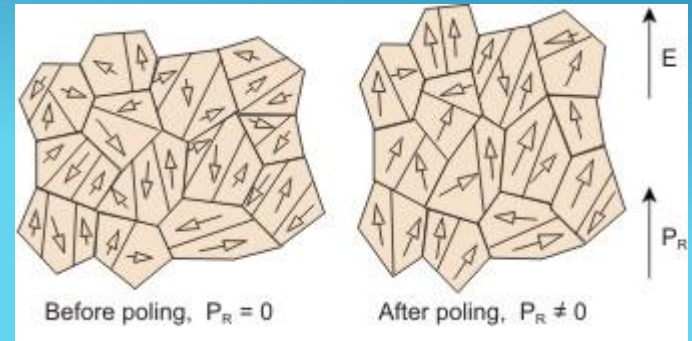


In BaTiO₃ at room temperature the stable phase has the Ti⁴⁺ ion displaced from the center of the cell. It has two stable positions: above and below the 4 central oxygens. Depending on its position, the dipole will be reversed.

Name (Abbreviation)	Chemical Formula	Curie Temperature, T_c (°C)	Spontaneous Polarization P_s ($\mu\text{C m}^{-2}$) at [T (°C)]	Crystal Structure	
				Above T_c	Below T_c
Barium Titanate	BaTiO ₃	120	26.0 [23]	Cubic	Tetragonal
Lead Titanate	PbTiO ₃	490	50.0 [23]	Cubic	Tetragonal
Potassium Niobate	KNbO ₃	435	30.0 [250]	Cubic	Tetragonal
Potassium Dihydrogen Phosphate (KDP)	KH ₂ PO ₄	-150	4.8 [-177]	Tetragonal	Orthorhombic
Triglycine Sulfate (TGS)	(NH ₂ CH ₂ COOH) ₃ • H ₂ SO ₄	49	2.8 [20]	Monoclinic (Centrosymm.)	Monoclinic (Noncentrosymm.)
Potassium-Sodium Tartrate-Tetrahydrate (Rochelle salt)	KNaC ₄ H ₄ O ₆ •4H ₂ O	24	0.25 [5]	Orthorhombic (Centrosymm.)	Monoclinic (Noncentrosymm.)
Antimony Sulfo-iodide	SbSI	22	25.0 [0]	Orthorhombic (Centrosymm.)	Orthorhombic (Noncentrosymm.)
Guanidinium Aluminium Sulfate Hexahydrate (GASH)	C(NH ₂) ₃ Al(SO ₄) ₂ •6H ₂ O	None	0.35 [23]	Trogonal	—

POLING

- Under a strong external field it is possible to induce the poling of the ferroelectric material (i.e. to polarize it)

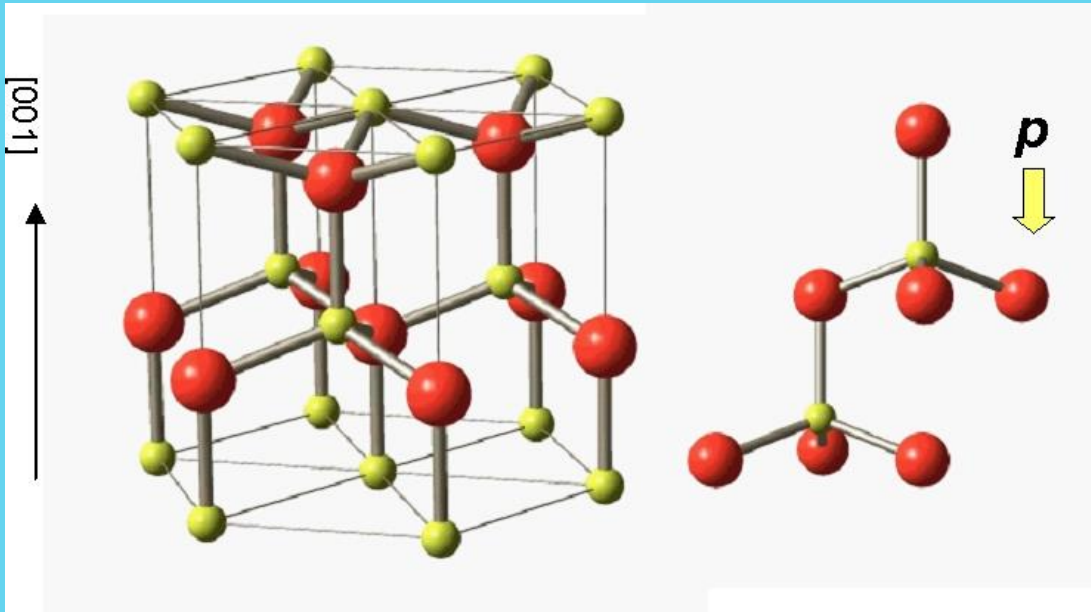


If the material is poled, the stress acting on it can generate an intense change of the electric dipole (usually stronger than in common piezoelectric materials)

PYROELECTRICS

Crystal with a permanent dipole, not reversible.

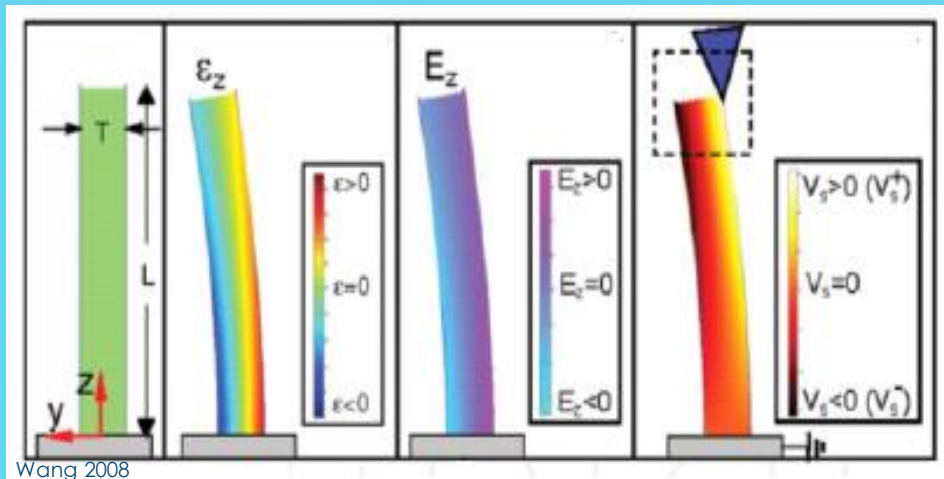
They have to be grown as single crystals.



Wurtzite structure crystals, such as ZnO, are asymmetric along the $[001]$ axis, ($[001]$ is different from $[00\bar{1}]$)

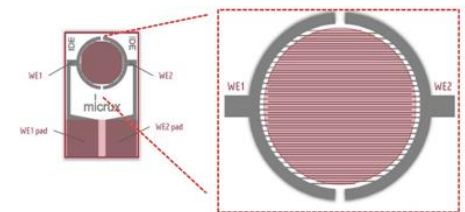
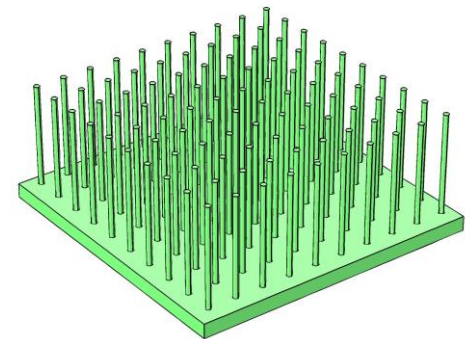
Because of thermal dilatation, the electric dipole increases

ZINC-OXIDE MICRORODS



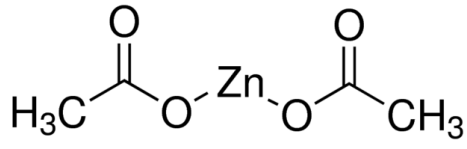
Objective: exploiting the difference of potential at the base of the pillar induced by the bending

Growth of ZnO pillars on IDE



FABRICATION

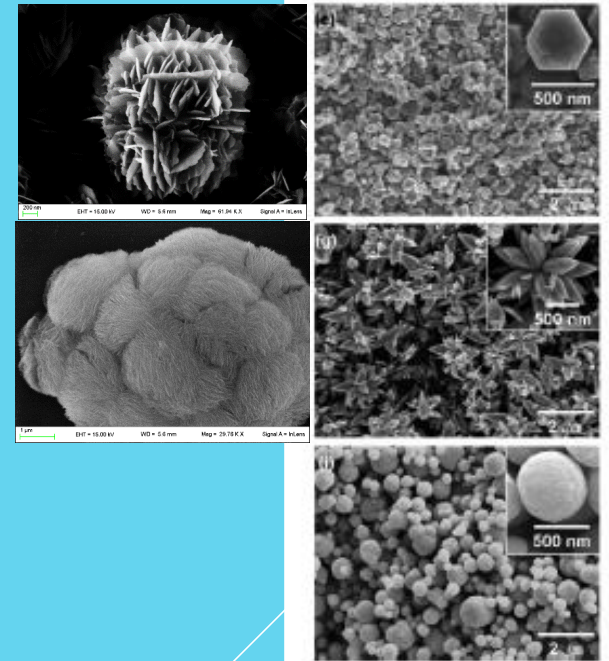
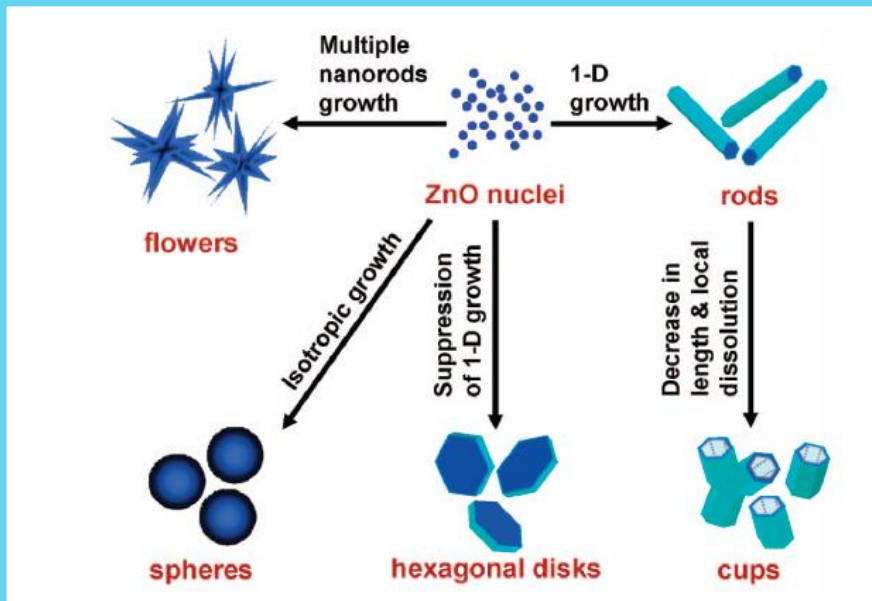
Hydrothermal Synthesis



Acetato di zinco

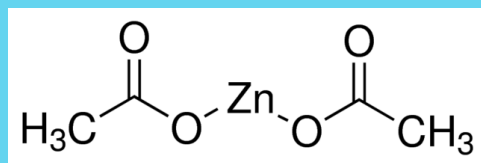
basic
environment

ZnO



HYDROTHERMAL SYNTHESIS

ZnO nano and microrods

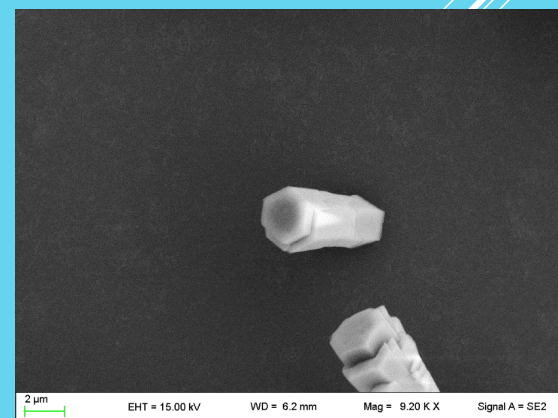
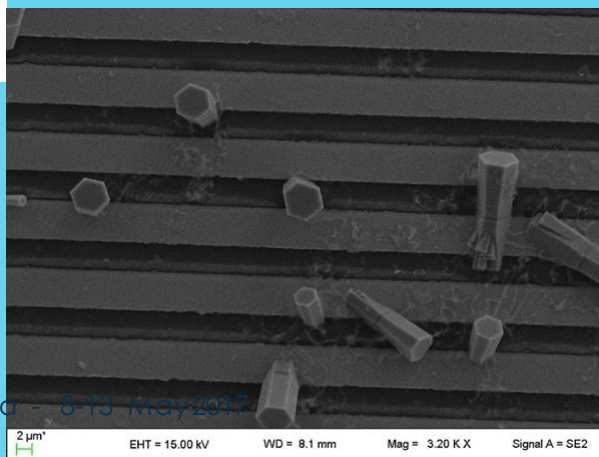
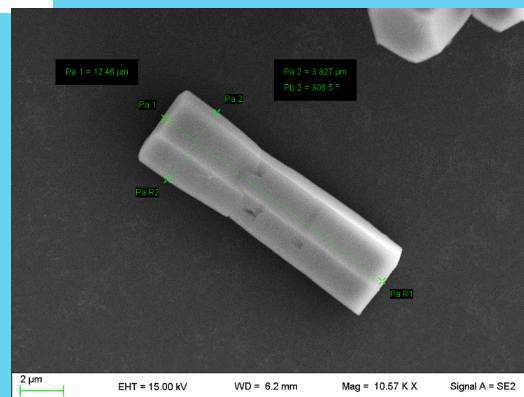
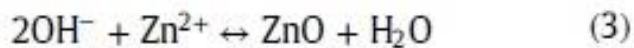
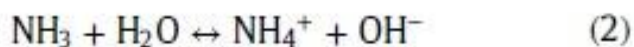
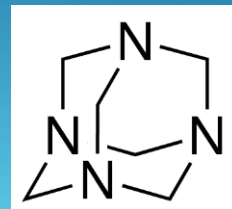


Acetato di zinco

HMTA

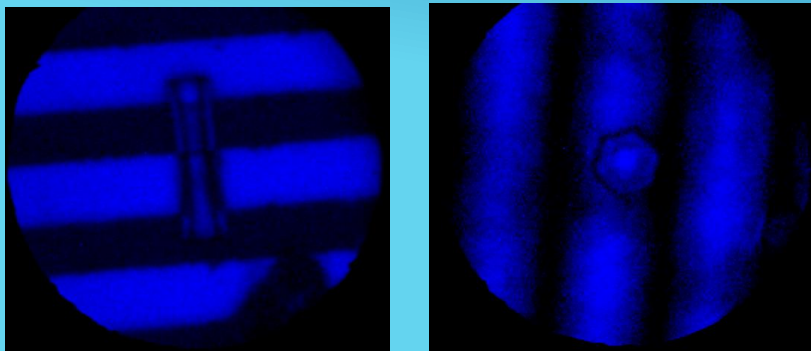
Hexamethylenetetramine

ZnO

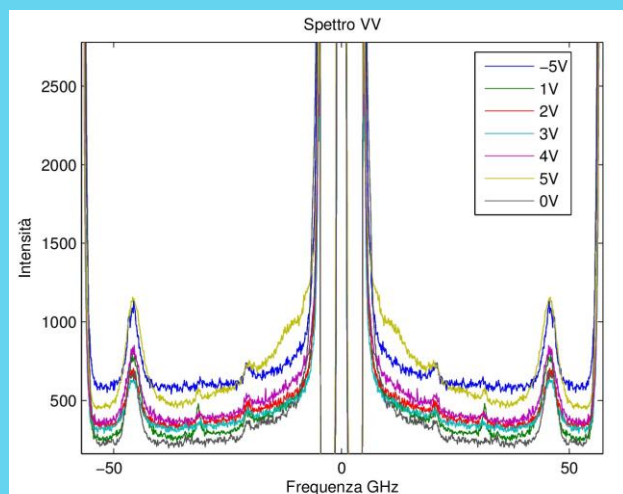


Size in the order of tens of micrometers

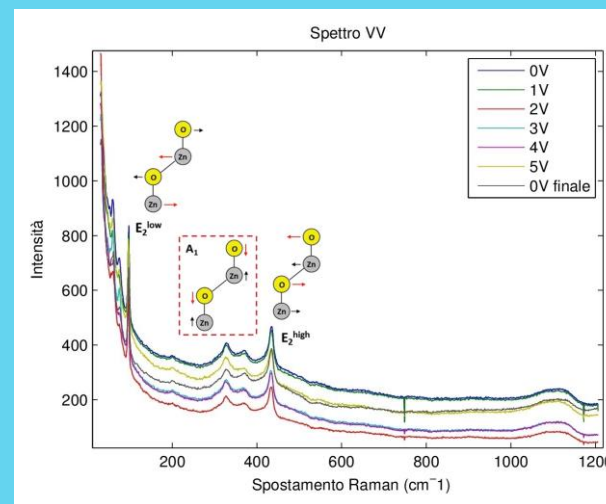
SPECTROSCOPIC CHARACTERIZATION



Raman and Brillouin spectroscopy on single crystals

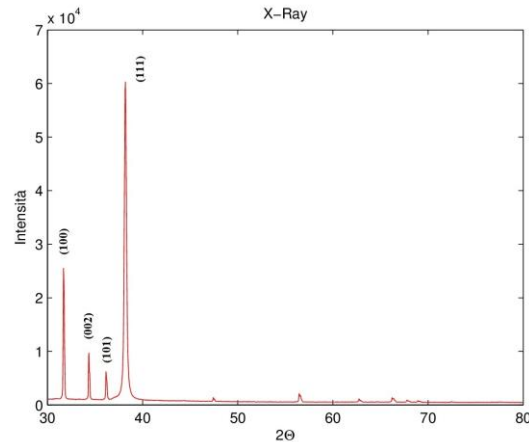
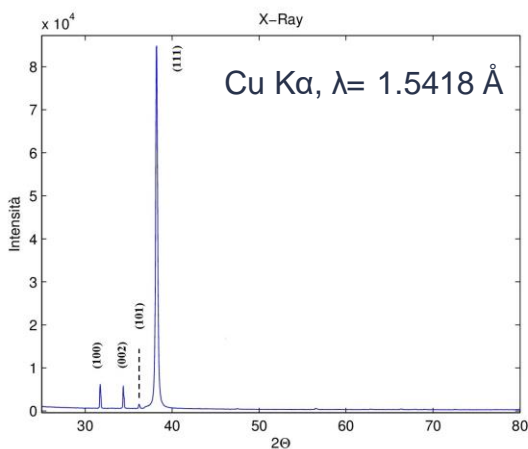


Elementi Matrice C (GPa)	C_{11}	C_{33}	C_{44}	C_{66}
Ref. Bhat et al.	209	210	42	44
Risultati	209	198	42	43



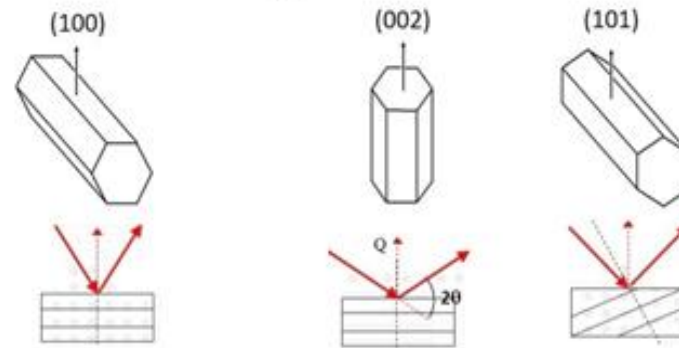
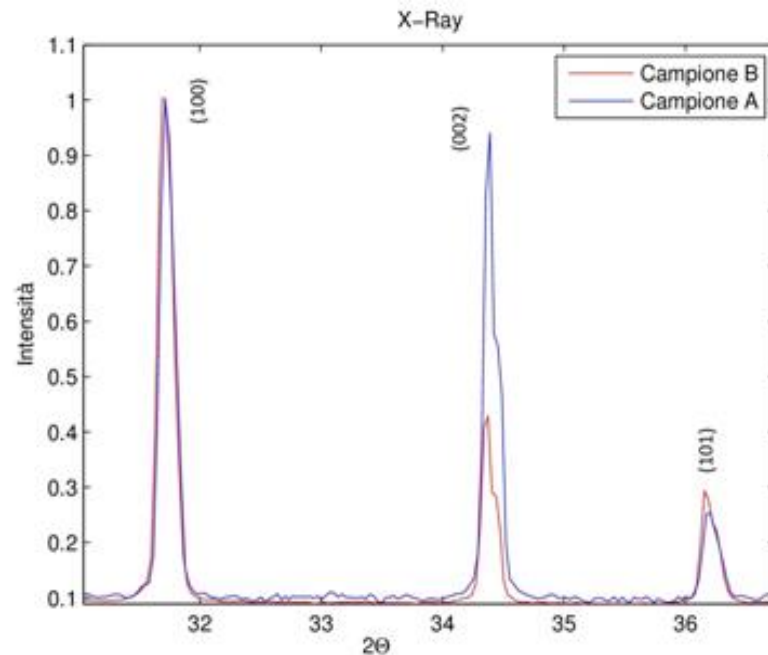
Raman Shift (cm⁻¹)	E_2^{low}	$A_1(E_2)$	$A_1(E_1,E_2)$	A_1	E_2^{high}
Ref. Damen et al.	101	208	332	380	437
Risultati	101	205	332	379	438

XRD CHARACTERIZATION

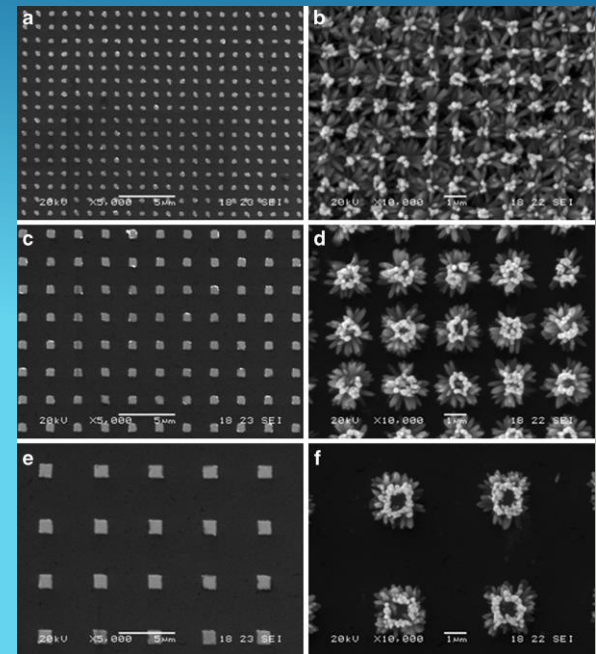
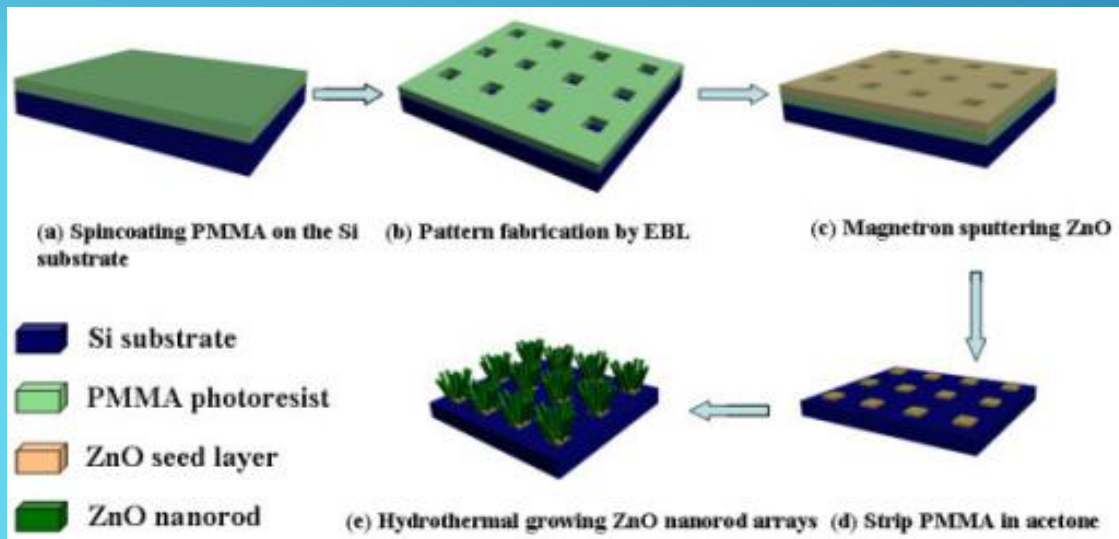


A (HMTA:ZnAc, 2:1)

B (HMTA:ZnAc, 3:1)

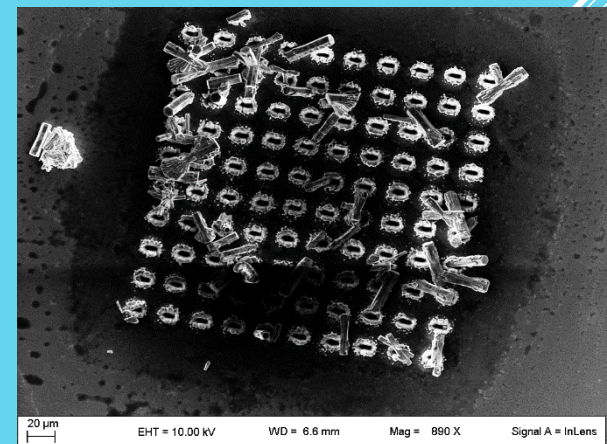
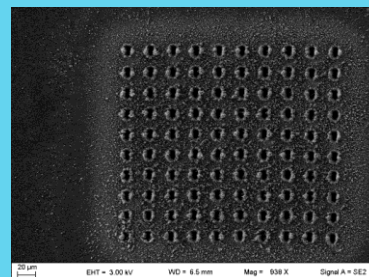


2θ (Deg)	(100)	(002)	(101)	(111)
Ref. Martinez et al.	31.73	34.40	36.21	38.19
Results	31.73	34.40	36.21	38.21



Wang et al. Nanoscale Research Letters 2012, 7:246

Template made by laser ablation on PMMA



ELECTRETS

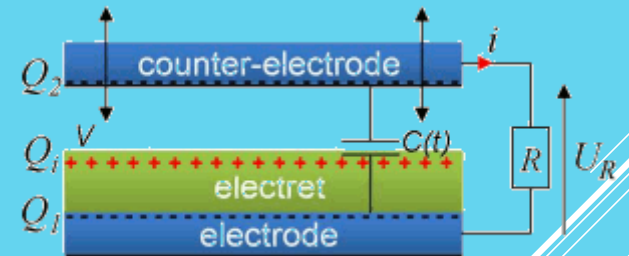
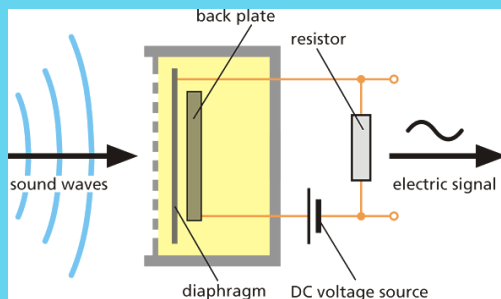
Dielectrics with unbalanced charge (permanent oriented electric dipoles or a net charge.)



Electrical analogue of a magnet: able to generate an electric field

MATERIALS: dielectrics (polymers, oxides) with high dielectric strength and low conductivity

Old applications: microphone



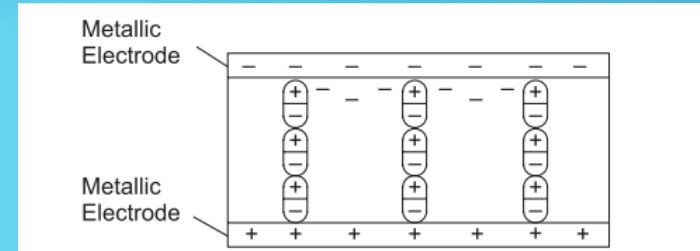
S Boisseau et al. Smart Materials and Structures 20, 105013, 2011

New ones : energy harvesting devices

FORMATION OF ELECTRETS

Two types of electrical charges in an electret :

- ▶ monocharges (also called real charges)
- ▶ dipolar charges (such as in ferroelectrets)

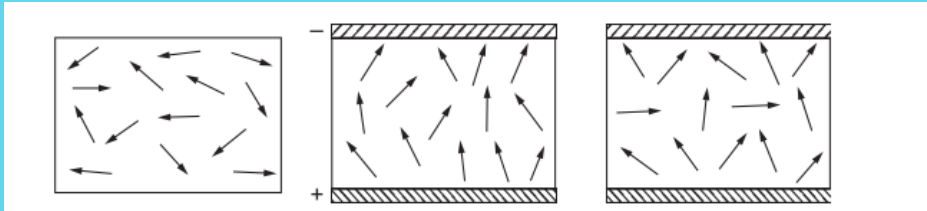


Several fabrication techniques:

- Thermo-Electrical Method (dipolar)
- Electromagnetic Radiation Method (dipolar)
- Liquid-Contact Method (real charges)
- Corona Discharge Method (real charges)
- Electron-Beam Method (real charges)

THERMO-ELECTRICAL METHOD

Dipolar molecules are randomly arranged but they will actively orient under an electric field at a temperature higher than the glass transition temperature, T_g



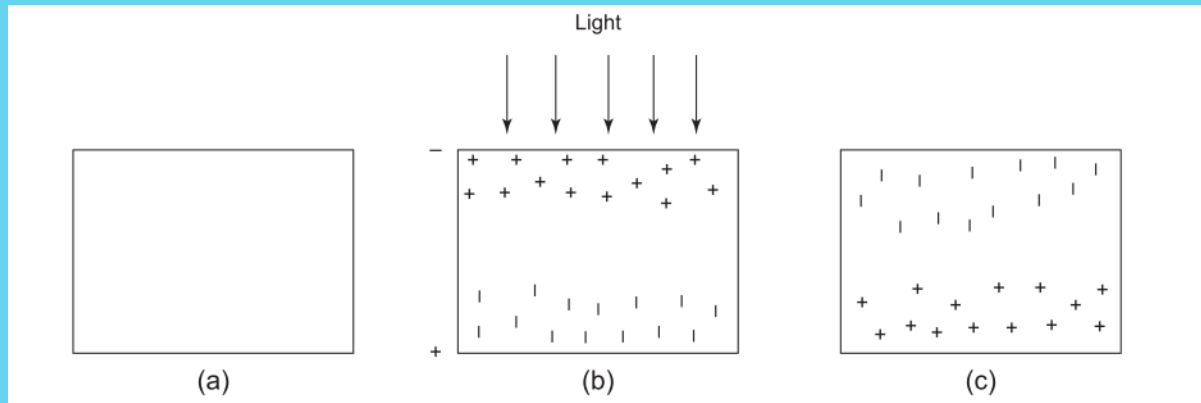
Typical materials are Carnauba wax–beeswax (first electret, made by Eguchi in 1919)

Drawback: stability

Ferroelectrets can be considered electrets obtained by thermoelectrical method (the poling) but with, possibly, much higher T_g

ELECTROMAGNETIC RADIATION METHOD

Displacement of the charge carriers generated by penetrating radiation (x-rays or ultraviolet light), under an externally applied electric field.

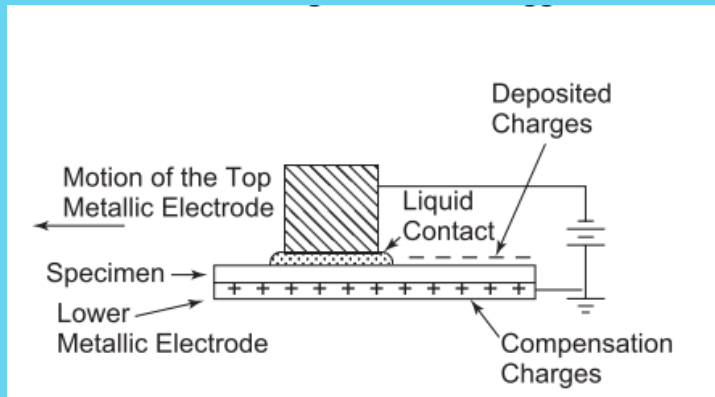


These carriers can be trapped near the electrodes to create a space charge polarization.

The polarization remains after the external field removal

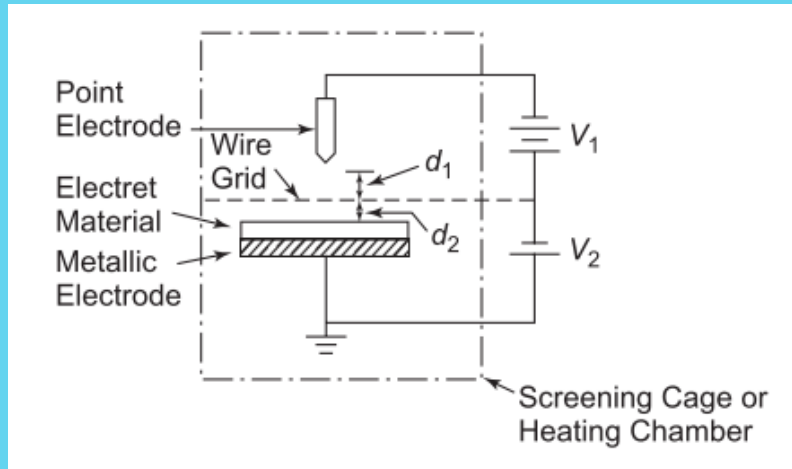
(LIQUID) CONTACT METHOD

Transfer of real charges into the material, by a conductive contact. This can be made at large scale down to the nanoscale (AFM)



The advantage of the conductive liquid is in the possibility to move the metallic contact all over a large surface

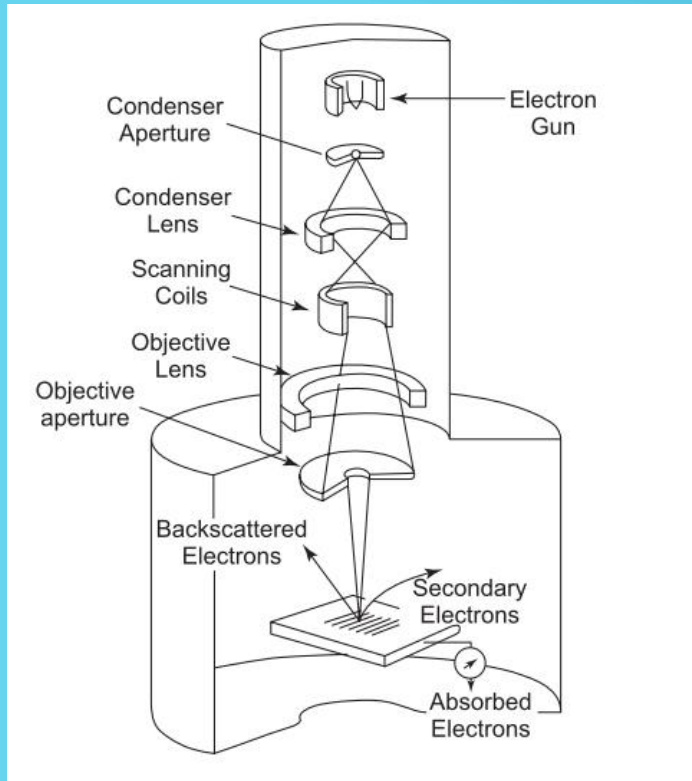
CORONA DISCHARGE



The bottom electrode is a vacuum-deposited metallic film on the material surface, and the top forming metallic electrode is usually made of a metallic wire

Around the point electrode it is possible to exceed the breakdown strength of the air in a region of a few millimeters. The so formed ions/free electrons can be accelerated toward the grid and so be implanted in the target dielectric material

ELECTRON BEAM METHOD



It is possible to inject/extract real charges (electrons) into the electret by SEM

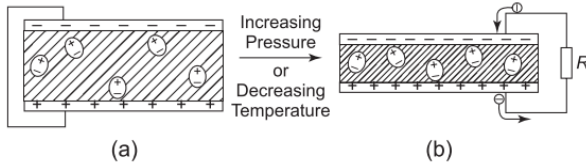
The energy of the electron beam (<50 keV) should be controlled according to the structure and thickness of the material specimens to be used for forming electret

A similar mechanism can be used with ion implantation instrument

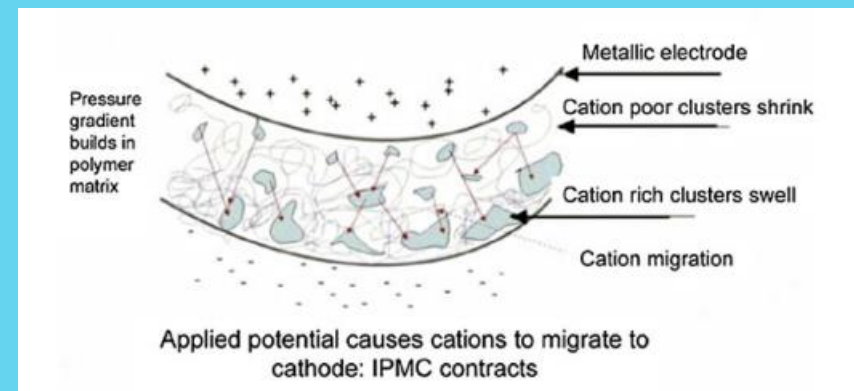
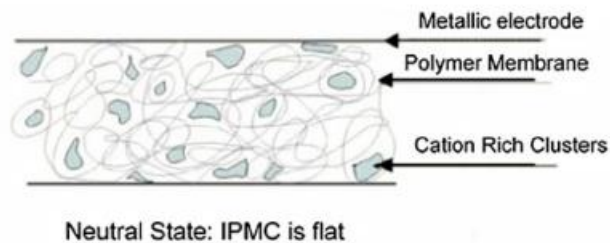
ELECTROACTIVE POLYMERS

(ARTIFICIAL MUSCLES)

Dipolar «soft» electrets have interesting properties similar to piezoelectric materials. They can react to external field changing shape or, viceversa, change their own dipolar field because of a change in shape.

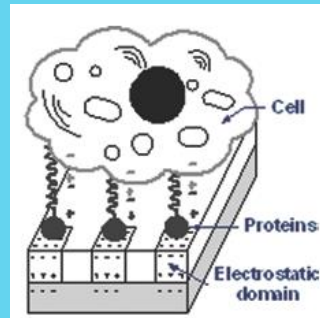
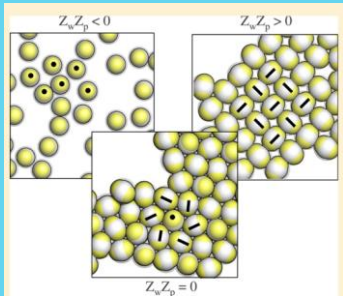


More complex structures can be engineered. Polymer are especially interesting.



SIZE REDUCTION OF ELECTRETS

- ▶ **Miniaturization** of devices (EH, MP)
- ▶ The charge provides a further way to **functionalize** the nano/micro- material



Tofail, Biological Interactions with Surface Charge in Biomaterials (RSC Publ.)

Bianchi et al, Nano Lett. 2014, 14, 3412 – 3418

Drawbacks

Stability of charging (surface vs space)
Control of charging

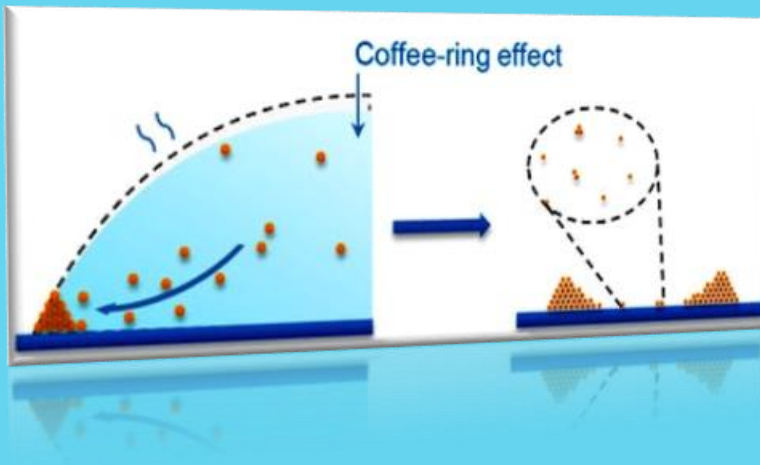
But also dynamic applications



Electro-Mechanical resonators

SAMPLES

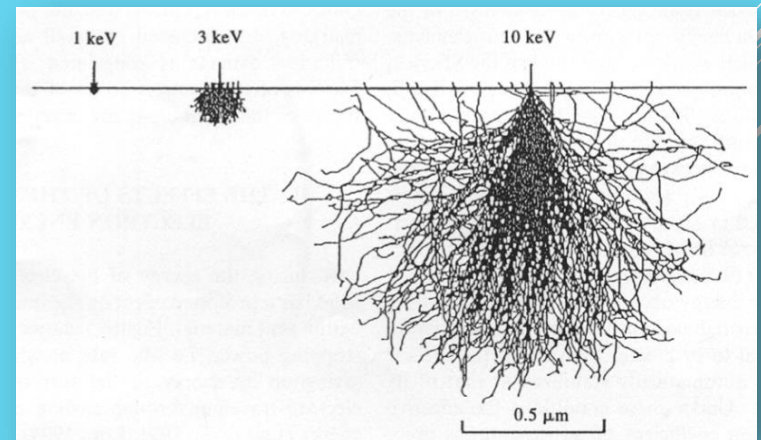
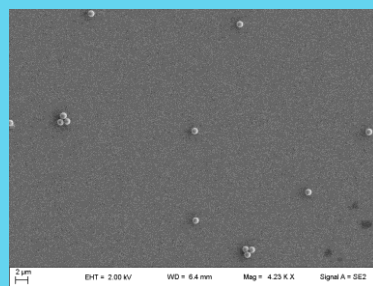
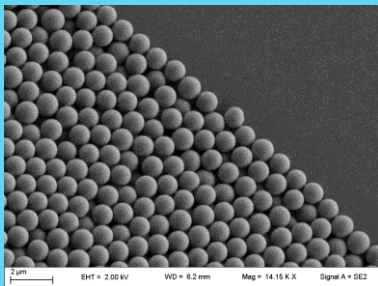
- ▶ SiO_2 particles (0.5 e 1 μm) deposited by drop casting on a polished copper substrate



Fabrication by SEM

- 1) High lateral resolution (5 nm at 20 keV)
- 2) Energy dependent penetration

$$R = (76/\rho)E_0^{1.67}$$



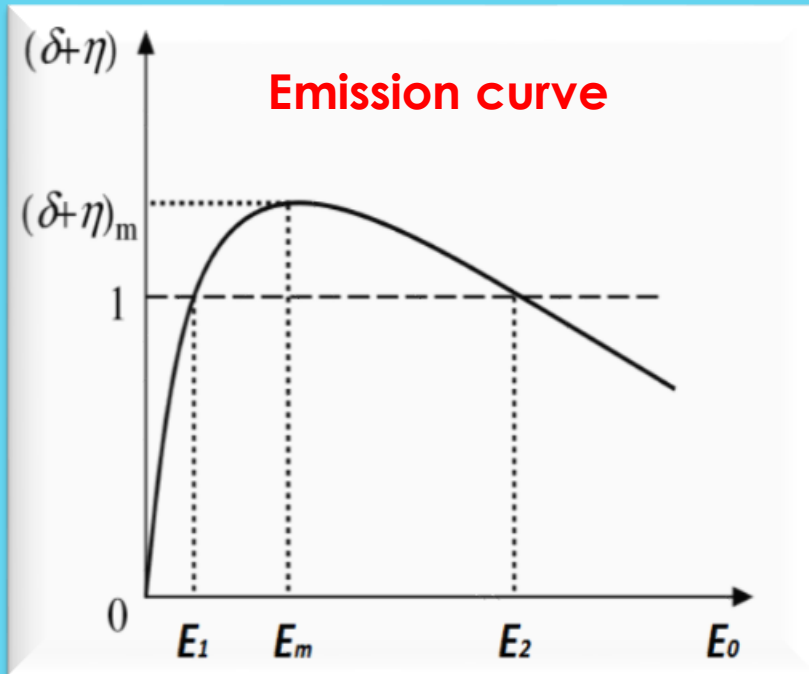
CHARGING MECHANISM

Total Yield Approach

$$\frac{\partial Q}{\partial t} = (1 - \sigma(E_0)) \cdot I_B - I_L$$

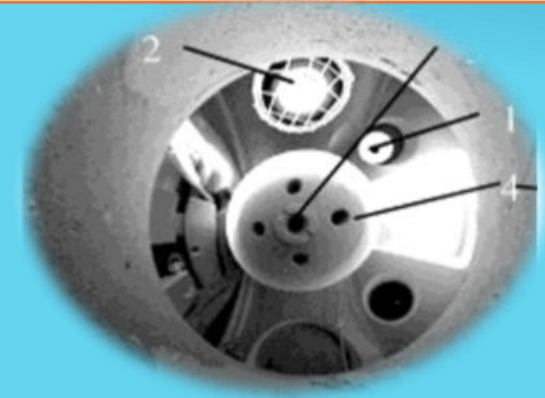
1) Crossover energies (E_1 e E_2)

2) Three charging region: ($\sigma > 1$ e $\sigma < 1$)



Charge effect:

$$E_L = E_0 - eV_s$$



Small size material:

1) Leakage current I_L can be significant

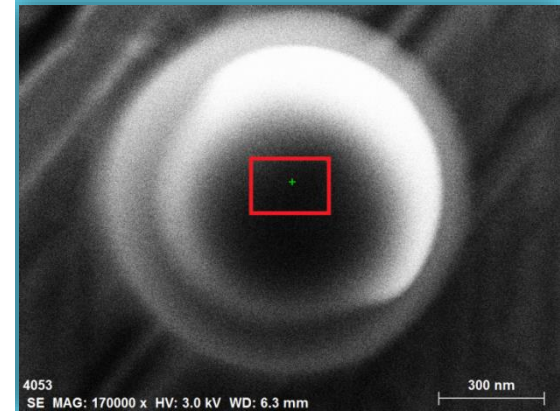
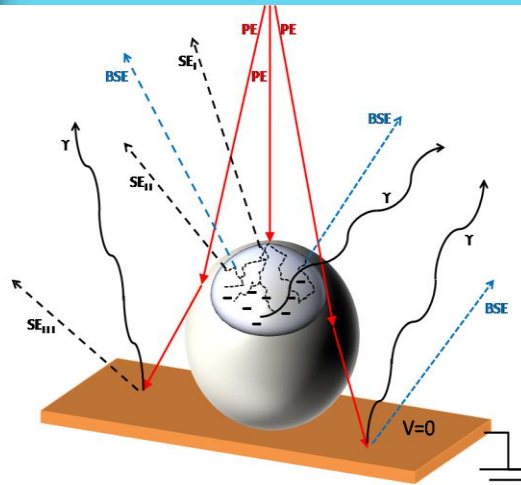
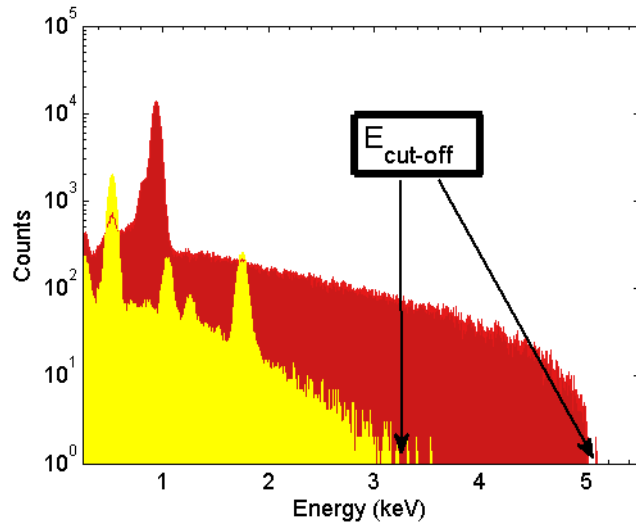
2) Surface potential V_s up to 200-300 V

DUANE-HUNT LIMIT SHIFT

$$E_L = E_0 - eV_s$$

Bremsstrahlung x-ray spectrum

$$V_s = \frac{E_0 - E_{cut-off}}{e}$$



$$\eta(\pi/2) \rightarrow 1$$

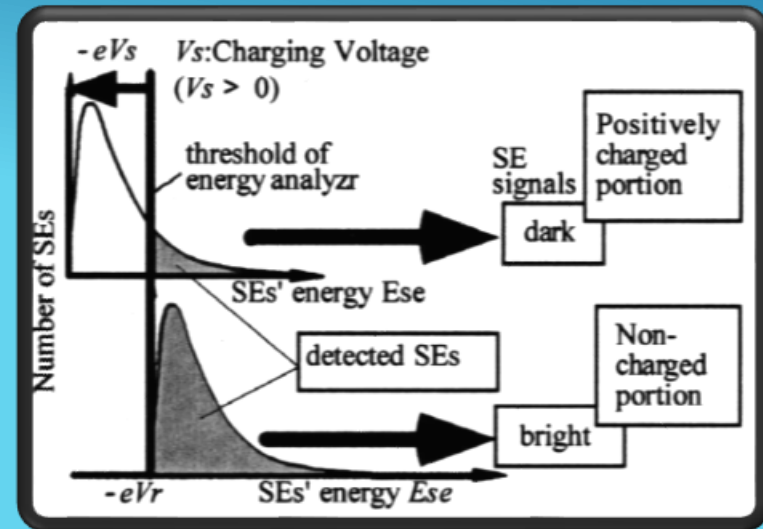
$$V_s < 2R \cdot r.d. \approx 350V \Rightarrow Q \approx 10^5 e^-$$

ELECTRONIC SPECTROSCOPY

Voltage Contrast

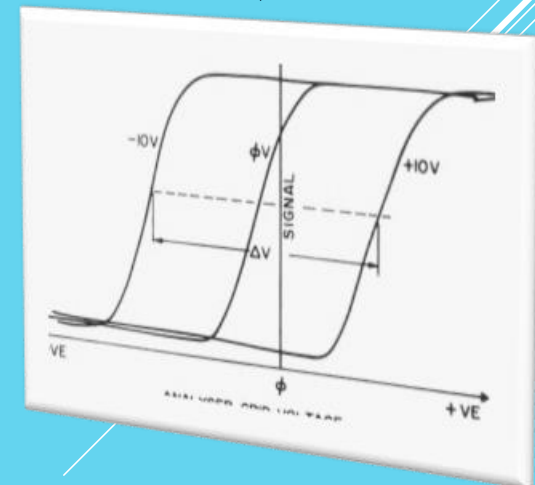
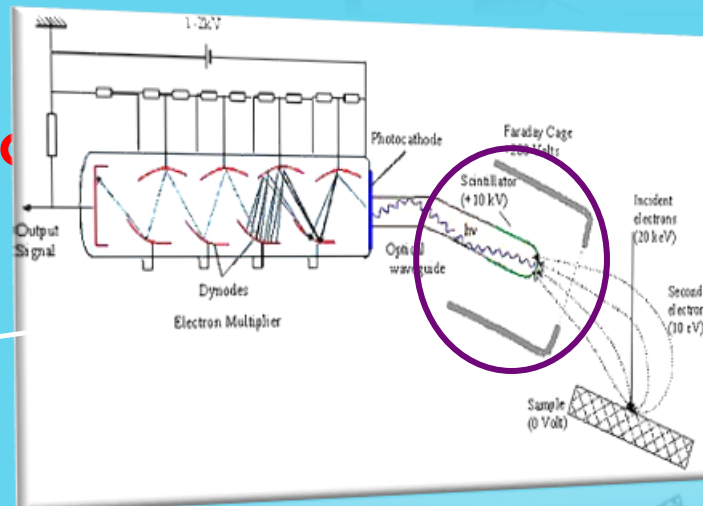
Chung-Everhart

$$\frac{dN}{dE} = \frac{k}{E_0} \frac{E}{(E + \phi)^4}$$

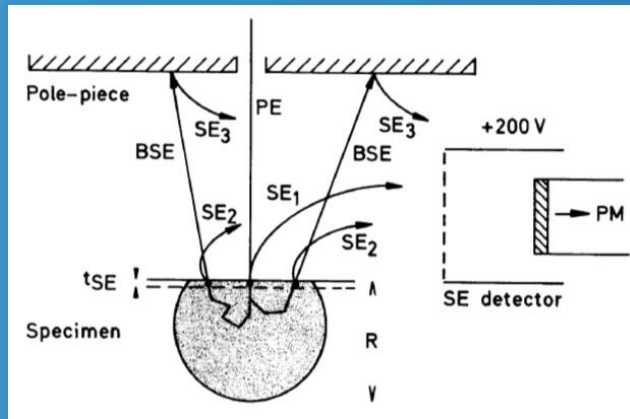


- Output depends on surface potential

Everhart-Thornley detector (E-T)



ELECTRONIC SPECTROSCOPY



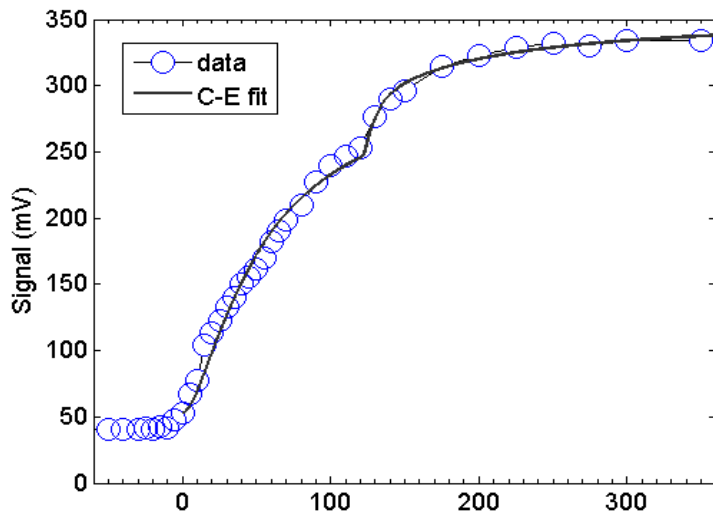
Strong background from the SEM chamber

Increasing Working Distance WD

Increasing Φ

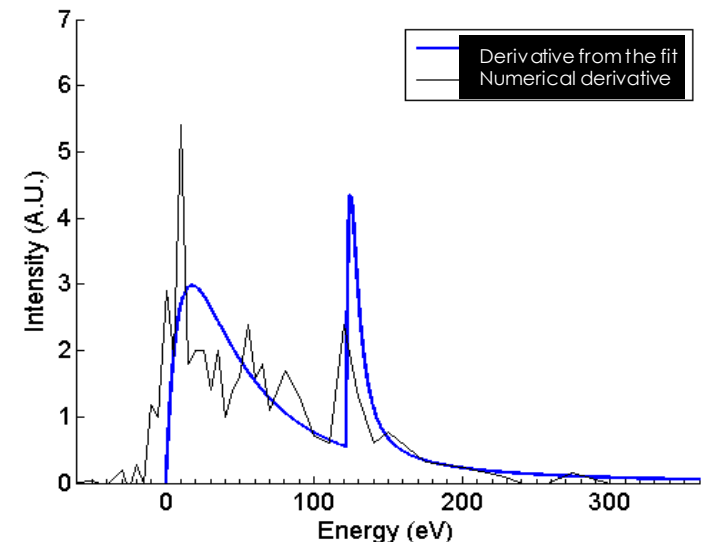
$$\frac{dN}{dE} = \frac{k}{E_0} \frac{E}{(E + \phi)^4}$$

$$S_{SE} = f_1(\delta \sec(\phi) + \delta \eta \beta) + f_2 \delta_{ext} + f_3 (d\eta / d\Omega) \Delta\Omega$$

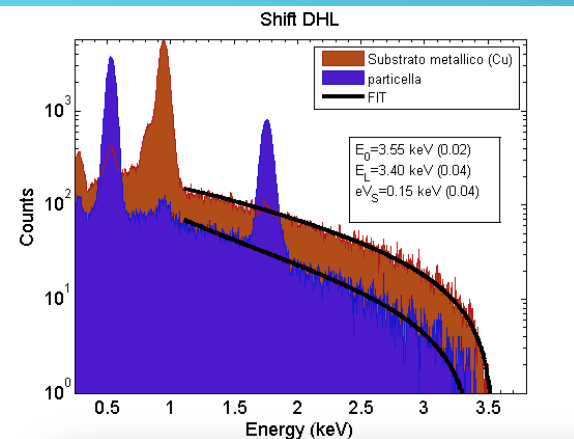


$E = 1.8 \text{ keV}$
 $WD = 6.8 \text{ mm}$
 $Mag = 150KX$

$V = (121 \pm 19) \text{ V}$
 $V_{rms} = (130 \pm 50) \text{ V}$

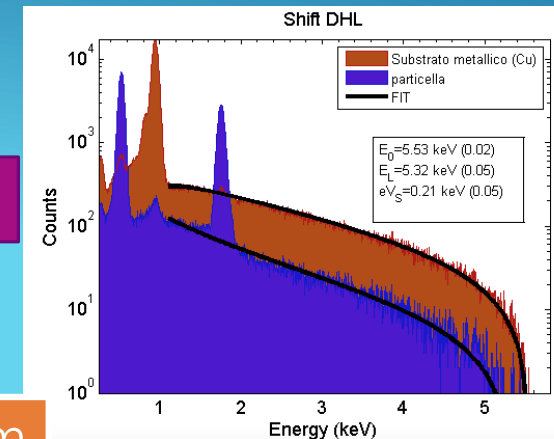
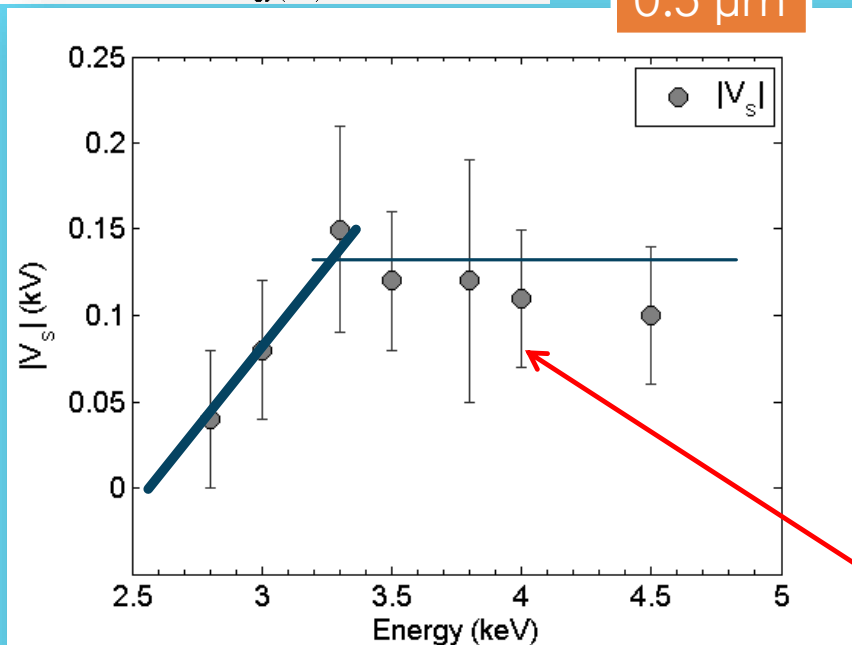


CHARGING VS ENERGY

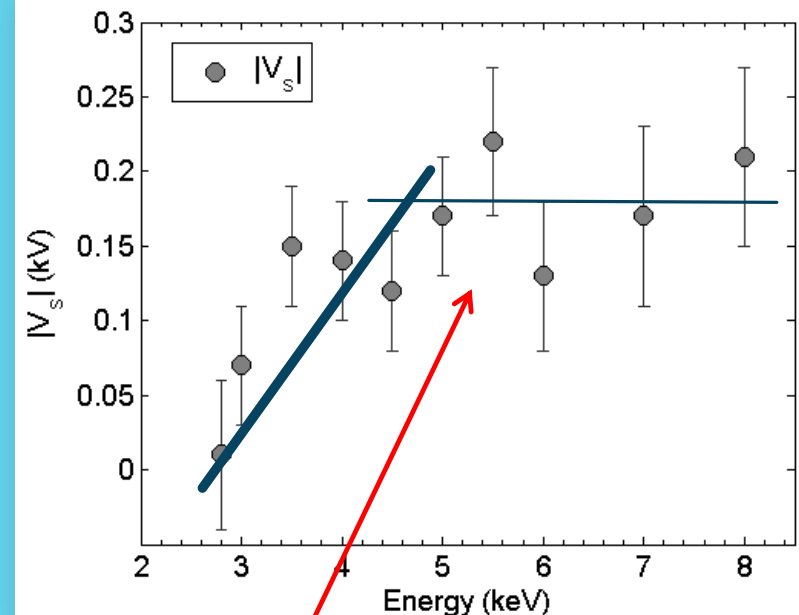


Initial linear increase

0.5 μm



1.0 μm



Potential saturation

CHARGE TEMPORAL EVOLUTION

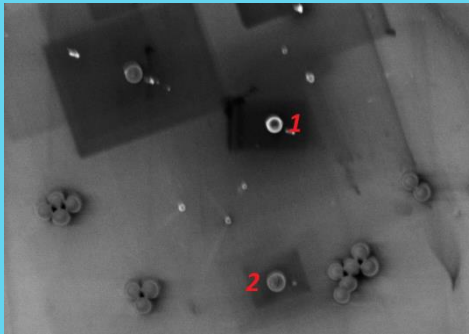
Samples in air

Monitoring charge by
non penetrating
electrons (0.5-1.5 keV)

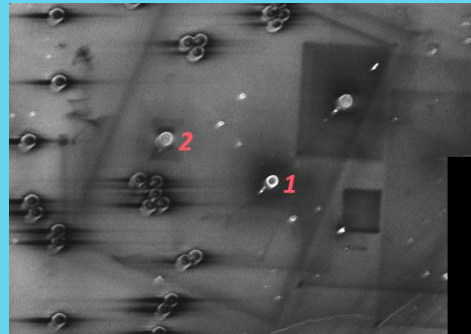
Observation
s by In Lens
detector



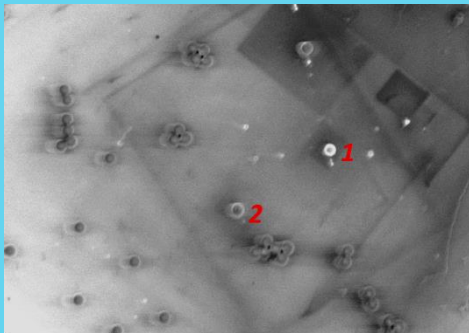
More electrons detected
from charged particles



13 days



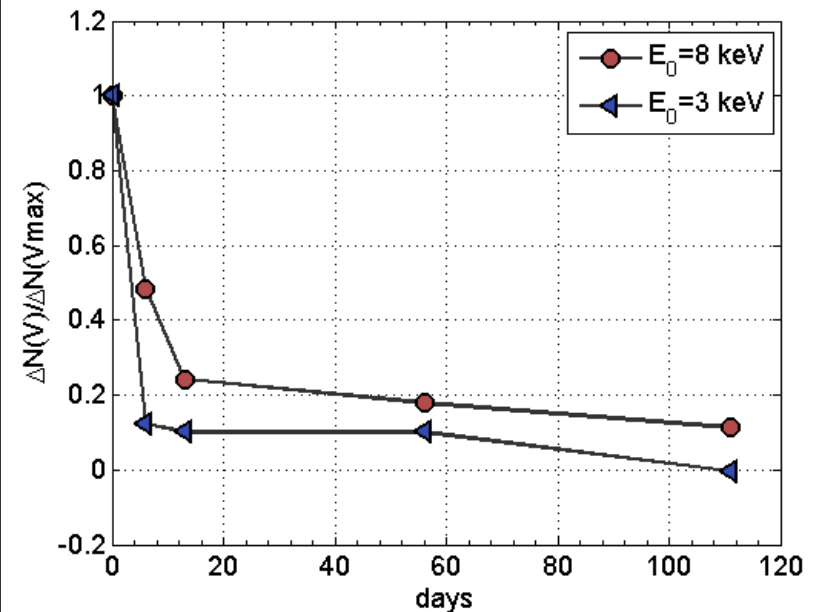
56 days



111 days

Two
characteristic
lifetimes:
1) **fast (surface)**
2) **Slow (space)**

$$\Delta N(V_s) = N_{charged}(V_s) - N_{neutral}$$



FINAL COMMENTS

- ▶ EH device and materials are strongly correlated
- ▶ Thermoelectric materials convert thermal gradient into electric energy. The optimization of thermoelectric generators depends on the separation of phononic and electronic properties.
- ▶ Electro active materials are effective ways to harvest mechanical energy (noise vibration and direct forces)
- ▶ Piezoelectrics materials change their polarization state after a shape change. Their properties depend on the asymmetric structure of the crystal cell.
- ▶ Electrets are artificial materials, where the natural charge distribution has been modified. They can provide significant permanent external electric fields.